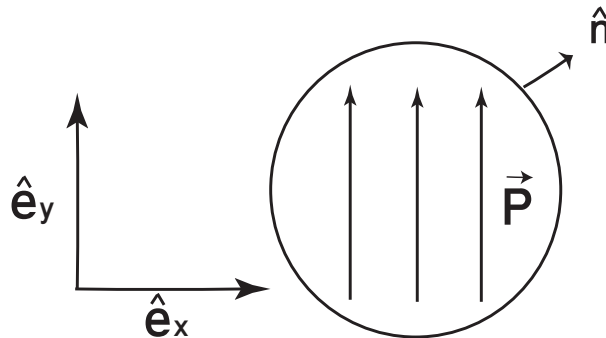


Solution HW 6: Dielectrics and Polarization

1. Answer:



Inside the sphere, $\rho_p = -\nabla \cdot \mathbf{P} = 0$. Since $\mathbf{P} = \text{constant}$. On the sphere surface $\sigma_p(\mathbf{r}_s) = \mathbf{P} \cdot \hat{\mathbf{n}}$, where

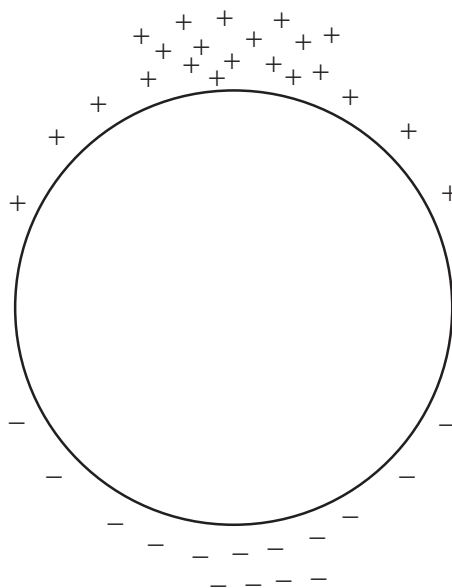
$$\hat{\mathbf{n}} = \sin \theta \hat{\mathbf{e}}_x + \cos \theta \hat{\mathbf{e}}_z,$$

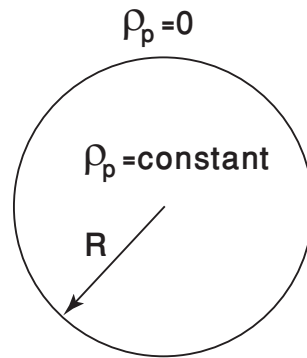
and θ is the angle with the vertical axis. Since

$$\mathbf{P} = P \hat{\mathbf{e}}_z$$

$$\begin{aligned} \sigma_p(\mathbf{r}_s) &= P \hat{\mathbf{e}}_z \cdot (\sin \theta \hat{\mathbf{e}}_x + \cos \theta \hat{\mathbf{e}}_z) \\ &= P \cos \theta. \end{aligned}$$

The varying charge on the surface of the sphere can be depicted as below.



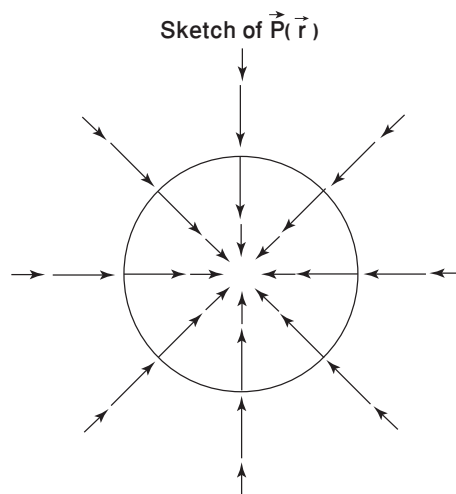
2. Answer:

Gauss's law

$$\nabla \cdot \mathbf{E} = +\frac{\rho}{\varepsilon_0} \quad (\rho \text{ inside a sphere})$$

has a solution:

$$\mathbf{E}(\mathbf{r}) = \begin{cases} \frac{\rho}{3\varepsilon_0} \mathbf{r}, & r < R \\ \frac{\rho R^3}{3\varepsilon_0 r^3} \mathbf{r}, & r > R \end{cases}$$



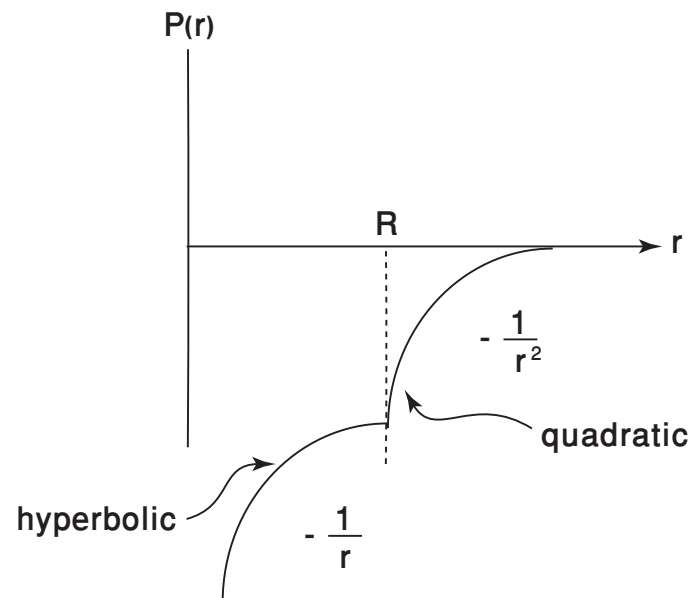
Similarly, here we need to solve,

$$\nabla \cdot \mathbf{P} = -\rho_p$$

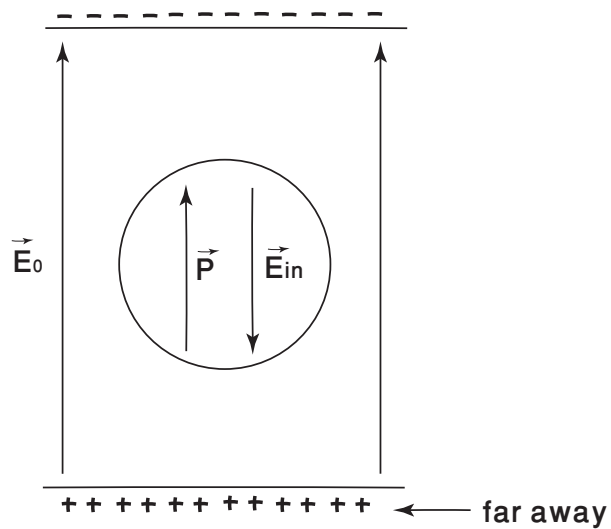
leading to (by analogy)

$$\mathbf{P}(\mathbf{r}) = \begin{cases} -\frac{\rho_p}{3} \mathbf{r}, & r < R \\ -\frac{\rho_p R^3}{3 r^3} \mathbf{r}, & r > R \end{cases}$$

\mathbf{P} is non-uniform and is non-zero outside the sphere too. A rough sketch is shown in the accompanying figure.



3. Answer:



$$\text{Now } \mathbf{E}_{\text{in}} = -\frac{\mathbf{P}}{3\epsilon_0} \quad (\text{derived in class, } \mathbf{P} \text{ is uniform inside the dielectric})$$

$$\mathbf{E} \text{ (total inside the dielectric)} = \mathbf{E}_p + \mathbf{E}_0$$

$$\begin{aligned} \text{Now } \mathbf{P} &= -3\epsilon_0 \mathbf{E}_p \\ &= -3\epsilon_0 (\mathbf{E} - \mathbf{E}_0) \end{aligned}$$

We also have

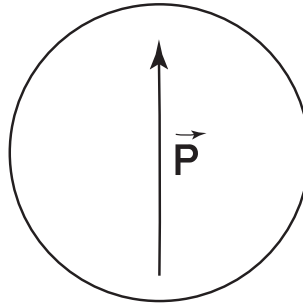
$$\begin{aligned}
 \mathbf{P} &= \varepsilon_0 \chi_e \mathbf{E} \\
 \therefore \varepsilon_0 \chi_e E &= -3\varepsilon_0(E - E_0) \\
 \varepsilon_0 \chi_e E + 3\varepsilon_0 E &= 3\varepsilon_0 \mathbf{E}_0 \\
 E(\chi_e + 3) &= 3E_0 \\
 E_0 &= \frac{E(\chi_e + 3)}{3} \\
 \text{Since } \chi_e &= \kappa - 1 \\
 E &= \left(\frac{3}{\kappa + 2} \right) E_0 \\
 \text{Hence } \mathbf{E} &= \mathbf{E}_0 \left(\frac{3}{\kappa + 2} \right)
 \end{aligned}$$

$$\begin{aligned}
 \text{Similarly } E_p &= E - E_0 \\
 &= E_0 \left(\frac{3}{\kappa + 2} - 1 \right) = E_0 \left(\frac{3 - \kappa - 2}{\kappa + 2} \right) = E_0 \left(\frac{1 - \kappa}{\kappa + 2} \right) \\
 \text{Hence } \mathbf{E}_p &= \mathbf{E}_0 \left(\frac{1 - \kappa}{\kappa + 2} \right)
 \end{aligned}$$

$$\text{Finally } \mathbf{P} = -3\varepsilon_0 \left(\frac{1 - \kappa}{\kappa + 2} \right) \mathbf{E}_0.$$

4. **Answer:**

(a)



$$\begin{aligned}\mathbf{E}_p(\text{in}) &= \mathbf{E}_{\text{in}} = -\frac{\mathbf{P}}{3\varepsilon_0}, & r < R \\ \mathbf{E}_p(\text{out}) &= \mathbf{E}_{\text{out}} = -\frac{V}{4\pi\varepsilon_0} \left(\frac{3(\mathbf{P} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{P}}{r^3} \right), & r > R\end{aligned}$$

Without loss of generality,

$$\mathbf{P} = P\hat{\mathbf{e}}_z.$$

θ is the polar angle from the vertical.

$$\mathbf{P} \cdot \hat{\mathbf{r}} = P \cos \theta$$

Inside the sphere

$$\begin{aligned}\mathbf{D}^{(\text{in})} &= \varepsilon_0 \mathbf{E}_{\text{in}} + \mathbf{P} = \varepsilon_0 \left(-\frac{\mathbf{P}}{3\varepsilon_0} \right) + \mathbf{P} = \frac{2}{3} \mathbf{P} \\ D_{\perp}^{(\text{in})} &= D_r^{(\text{in})} = \frac{2}{3} P \cos \theta.\end{aligned}$$

The normal component is the radial component.

Outside the sphere

$$\begin{aligned}\mathbf{P} &= 0 \\ \mathbf{D}^{(\text{out})} &= \varepsilon_0 \mathbf{E}_{\text{out}} = \frac{V}{4\pi} \left(\frac{3(\mathbf{P} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{P}}{r^3} \right) \\ D_{\perp}^{(\text{out})} &= D_r^{(\text{out})} = \frac{V}{4\pi r^3} (3P \cos \theta - P \cos \theta)\end{aligned}$$

If $r = R$

$$D_{\perp}^{(\text{out})} = D_r^{(\text{out})} = \frac{4}{3} \pi R^3 \frac{P}{4\pi R^3} (2 \cos \theta) = \frac{2P \cos \theta}{3}.$$

Hence D_{\perp} is continuous across the surface of the dielectric.

(b) We now compute the normal components of the electric field.

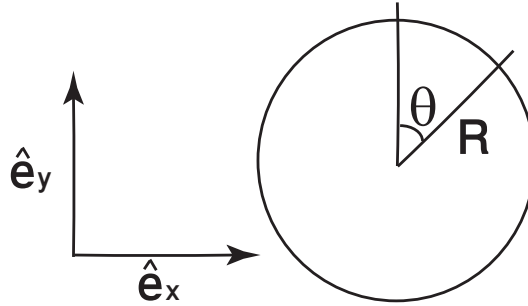
$$\begin{aligned}E_{\perp}^{(\text{in})} &= -\frac{P}{3\varepsilon_0} \cos \theta \\ E_{\perp}^{(\text{out})} &= \frac{2P}{3\varepsilon_0} \cos \theta\end{aligned}$$

Discontinuity is

$$E_{\perp}^{(\text{out})} - E_{\perp}^{(\text{in})} = \frac{P}{3\epsilon_0} \cos \theta (2 + 1) = \frac{P}{\epsilon_0} \cos \theta.$$

The electric field's \perp component is discontinuous, but the D 's \perp component is continuous.

(c) Find field at the center of a uniformly polarized sphere by explicit integration.



There is only charge on the surface (see Q1.).

Since $\sigma = P \cos \theta$

$$\begin{aligned} \mathbf{E}(\mathbf{r} = 0) &= \frac{1}{4\pi\epsilon_0} \int_S \frac{d^2 r' \sigma(\mathbf{r}')}{r'^3} (-\mathbf{r}) \\ &= -\frac{P}{4\pi\epsilon_0 R^3} \int_S d^2 r' \cos \theta \mathbf{R} \end{aligned}$$

Let's find the x -component:

$$\begin{aligned} E_x &= -\frac{P}{4\pi\epsilon_0 R^3} \int_S d^2 r' \cos \theta x \\ x &= R \sin \theta \cos \phi \\ E_x &= -\frac{P}{4\pi\epsilon_0 R^3} \int d\theta d\phi R^2 \sin \theta \cos \theta (R \sin \theta \cos \phi) \\ &= -\frac{P}{4\pi\epsilon_0} \left(\int_0^{2\pi} d\phi \cos \phi \right) \left(\int_0^\pi d\theta \sin^2 \theta \cos \theta \right) = 0 \end{aligned}$$

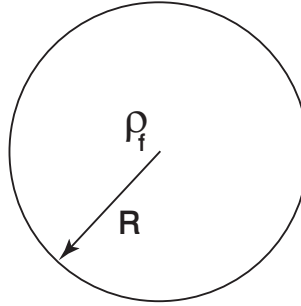
Similarly, by symmetry $E_y = 0$. whereas, for the z -component

$$\begin{aligned}
 E_z &= -\frac{P}{4\pi\epsilon_0 R^3} \int_s d^2r' \cos\theta z \\
 &= -\frac{P}{4\pi\epsilon_0 R^3} \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} d\theta d\phi R^2 \sin\theta \cos\theta (R \cos\theta) \\
 &= -\frac{P}{4\pi\epsilon_0} (2\pi) \int_0^{\pi} d\theta \sin\theta \cos^2\theta \\
 &= -\frac{P}{2\epsilon_0} \cdot \left(-\frac{\cos^3\theta}{3} \right) \Big|_0^{\pi} = \frac{P}{6\epsilon_0} (-2) = -\frac{P}{3\epsilon_0}.
 \end{aligned}$$

Hence \mathbf{E} at center is $-\frac{P}{3\epsilon_0}$ as desired.

5. **Answer:**

(a)



The density ρ_f is uniform. Let's find \mathbf{D} using Gauss's law:

$$\begin{aligned}
 \nabla \cdot \mathbf{D} &= \rho_f \\
 \int_S d\mathbf{S} \cdot \mathbf{D} &= D(4\pi r^2) = \frac{4}{3}\pi r^3 \rho_f \\
 \mathbf{D} &= \frac{r\rho_f}{3} \hat{\mathbf{r}} \\
 \text{Now } \mathbf{D} &= \epsilon_0 \mathbf{E} + \mathbf{P} \\
 \mathbf{P} &= \mathbf{D} - \epsilon_0 \mathbf{E} = \mathbf{D} - \frac{\epsilon_0}{\epsilon_0} \mathbf{D} = \mathbf{D} \left(1 - \frac{1}{\kappa} \right) \\
 \text{Hence } \mathbf{P} &= \frac{r\rho_f}{3} \hat{\mathbf{r}} \left(\frac{\kappa - 1}{\kappa} \right) = \mathbf{r} \frac{\rho_f}{3} \left(\frac{\kappa - 1}{\kappa} \right).
 \end{aligned}$$

(b)

$$\begin{aligned}\rho_p &= -\nabla \cdot \mathbf{P} = -\frac{\rho_f}{3} \left(\frac{\kappa - 1}{\kappa} \right) \nabla \cdot \mathbf{r} = -\rho_f \left(\frac{\kappa - 1}{\kappa} \right). \\ \sigma_p &= \mathbf{P} \cdot \hat{\mathbf{n}} = \frac{\rho_f}{3} \left(\frac{\kappa - 1}{\kappa} \right) \mathbf{r} \cdot \hat{\mathbf{n}} = \frac{\rho_f}{3} \left(\frac{\kappa - 1}{\kappa} \right) R.\end{aligned}$$

(c)

$$\text{Total polarization charge on surface} = \frac{\rho_f}{3} \left(\frac{\kappa - 1}{\kappa} \right) R (4\pi R^2) = Q_f \left(\frac{\kappa - 1}{\kappa} \right).$$