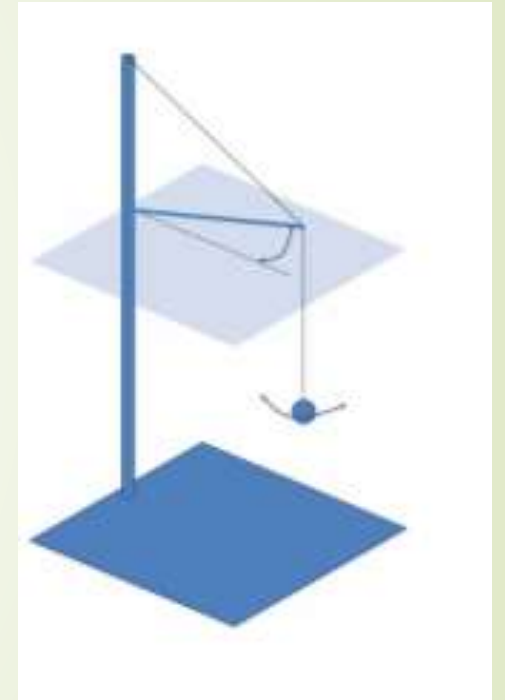


Problem 11. Azimuthal-Radial Pendulum

By: Aarij Atiq

Problem Statement

- Fix one end of a horizontal elastic rod to a rigid stand. Support the other end of the rod with a taut string to **avoid vertical deflection** and suspend a bob from it on another string. In the resulting pendulum the **radial oscillations** (parallel to the rod) can spontaneously **convert into azimuthal oscillations** (perpendicular to the rod) and vice versa. **Investigate the phenomenon.**



Outline of report

1

- Preliminary Observations

2

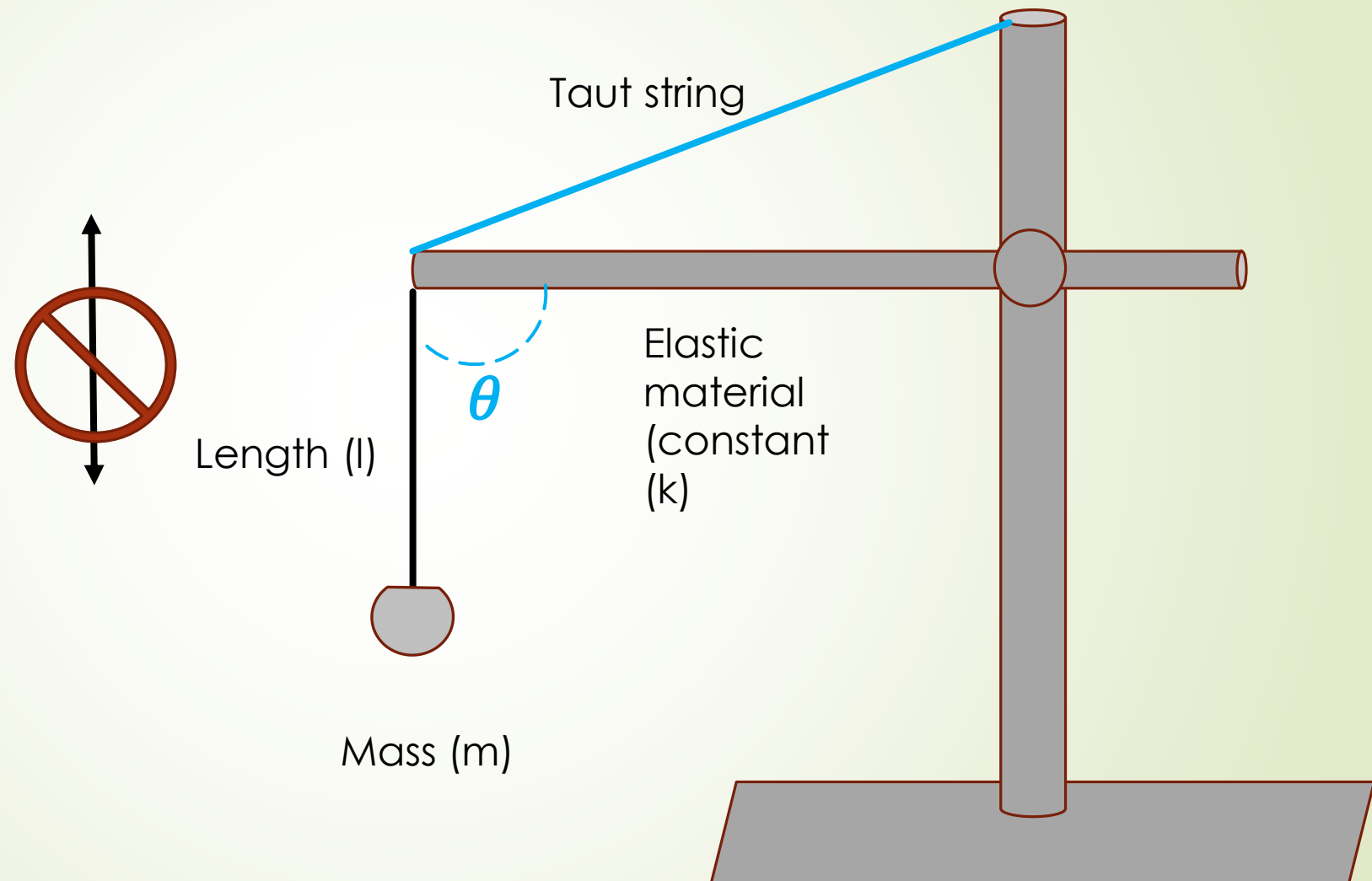
- Theory

3

- Experiments
- Experimental Set-up, and Results

- Conclusion

Construction of system



5

Diagram – Top View

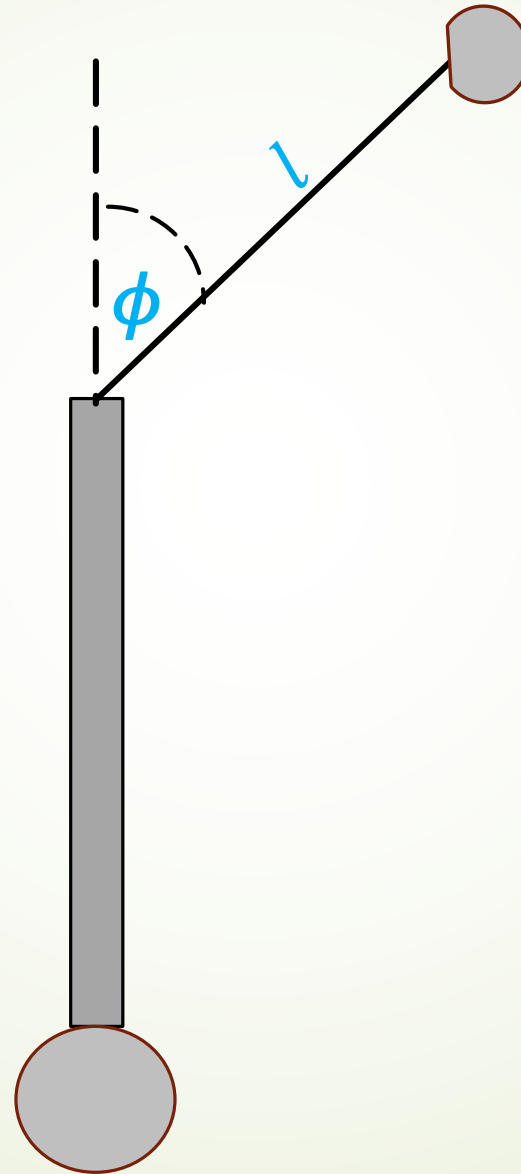
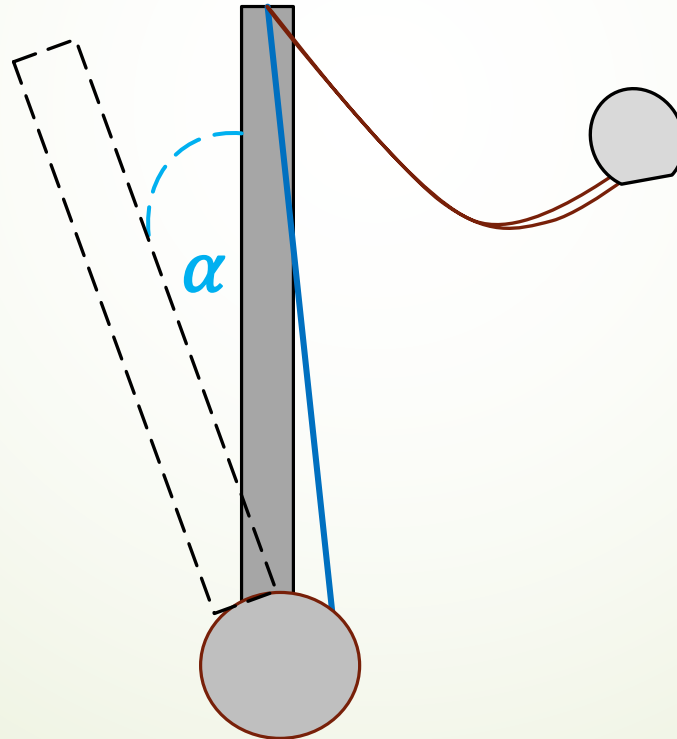


Diagram Top View



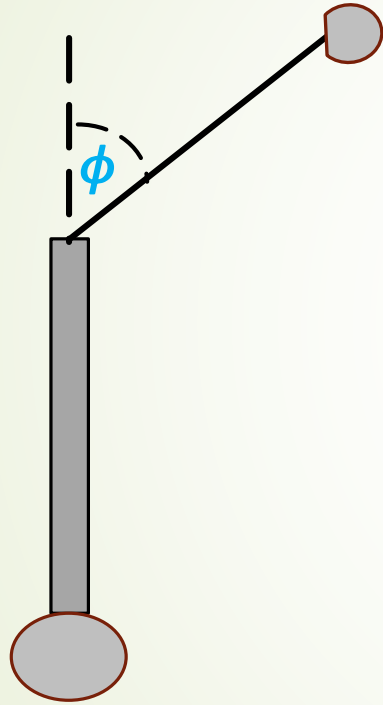
Observation Video

The pendulum behaves like a conical pendulum with an oscillating fulcrum

Radial oscillations → Azimuthal oscillations

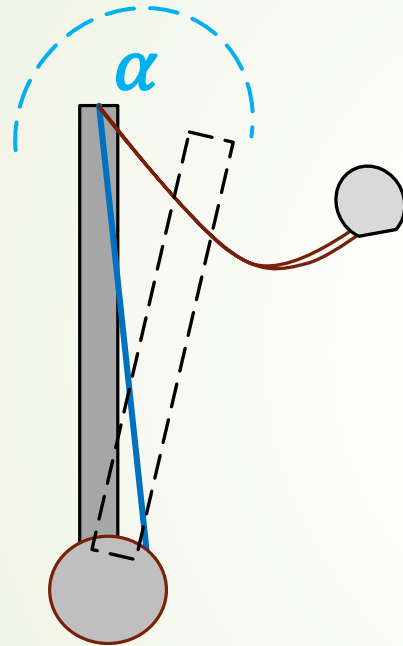
8

Investigation- Conditions of the Experiment



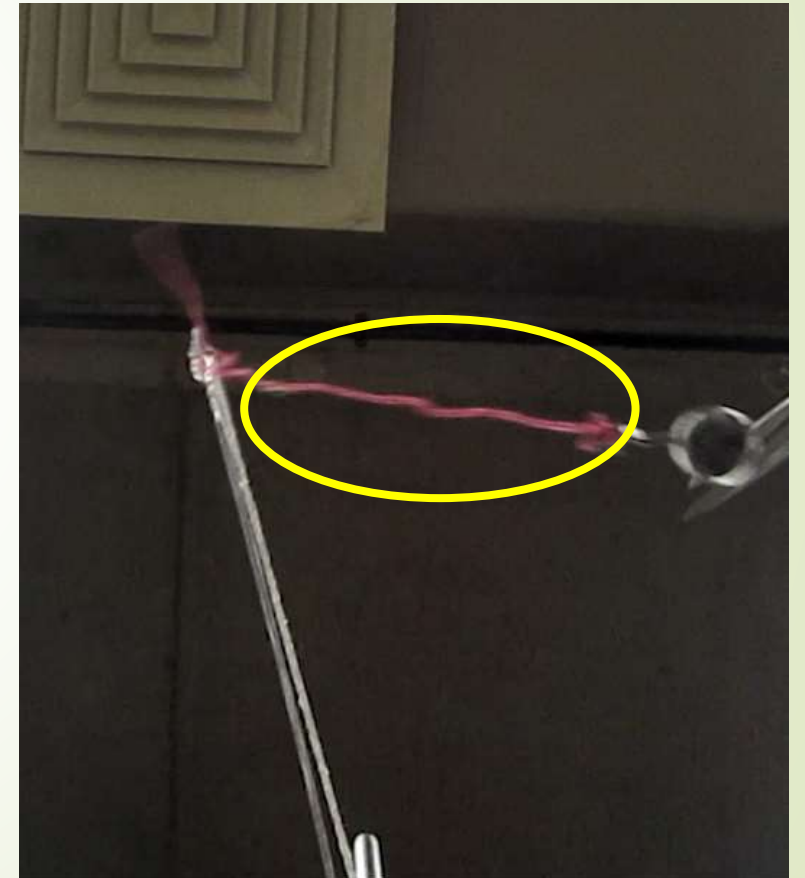
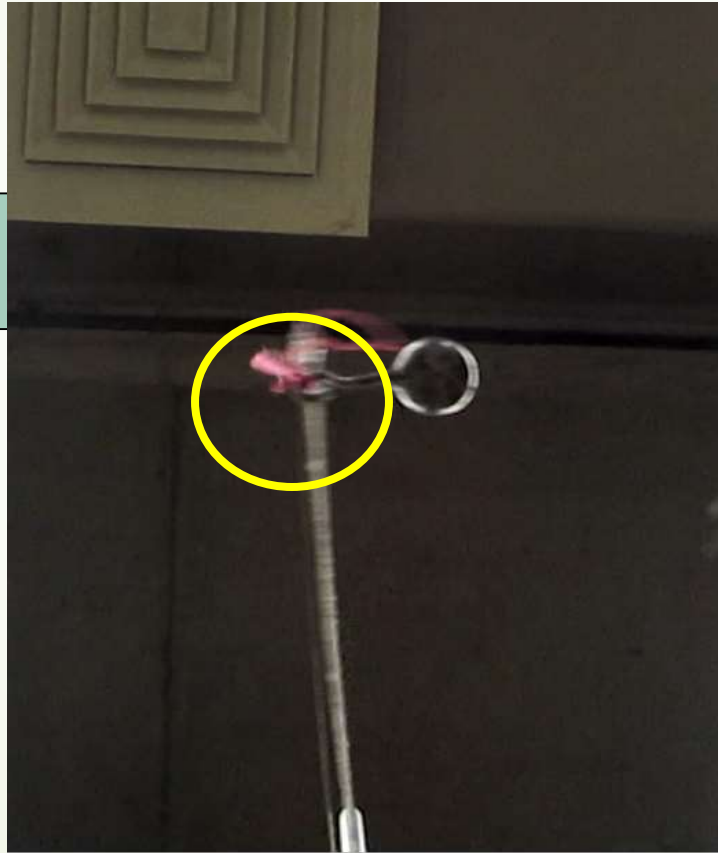
The pattern isn't observed when disturbed in the phi direction

Investigation-Conditions



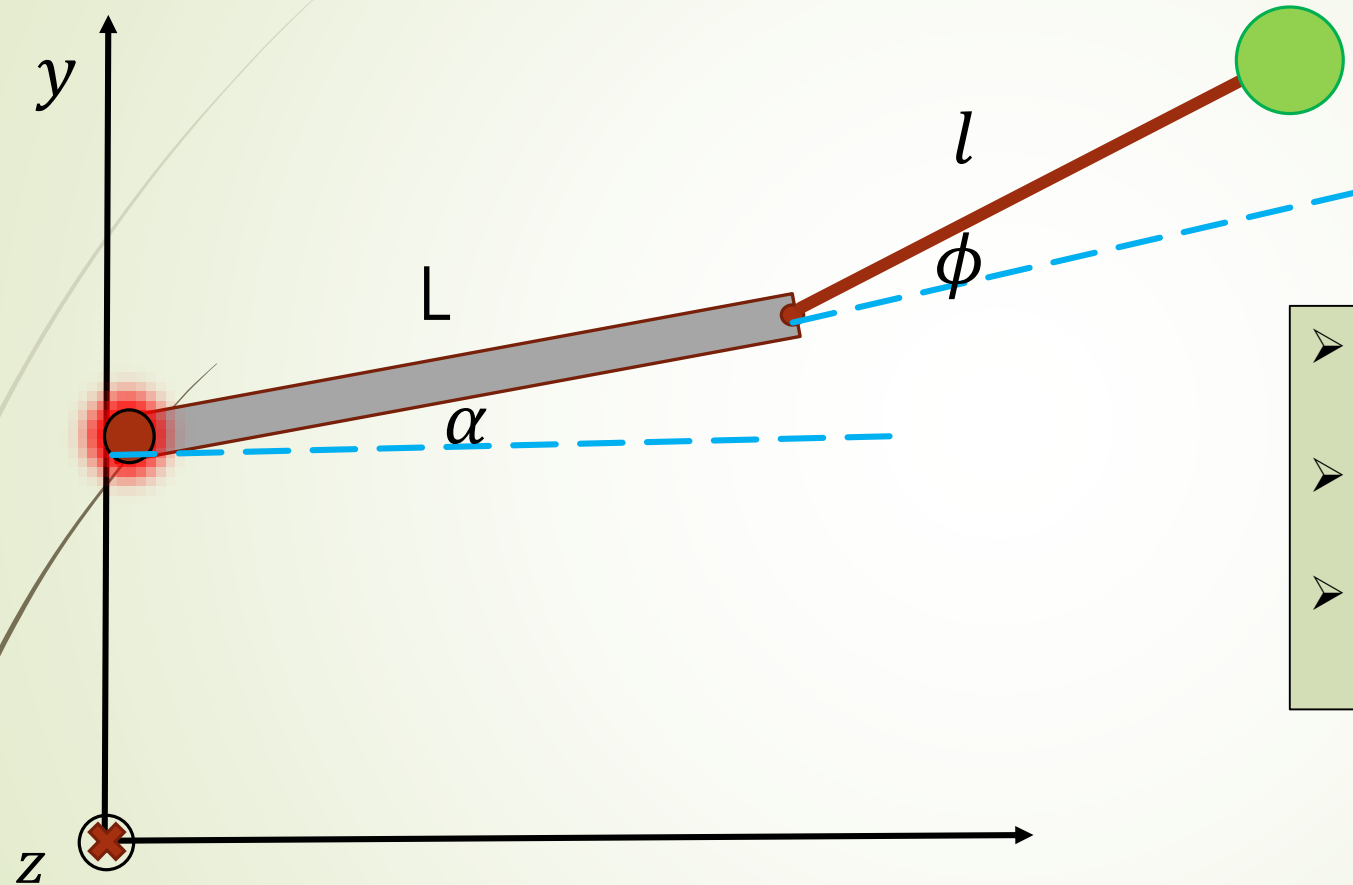
*Behaves like a simple
conical pendulum when
disturbed in the α direction*

Effect achieved due to the contraction of the string



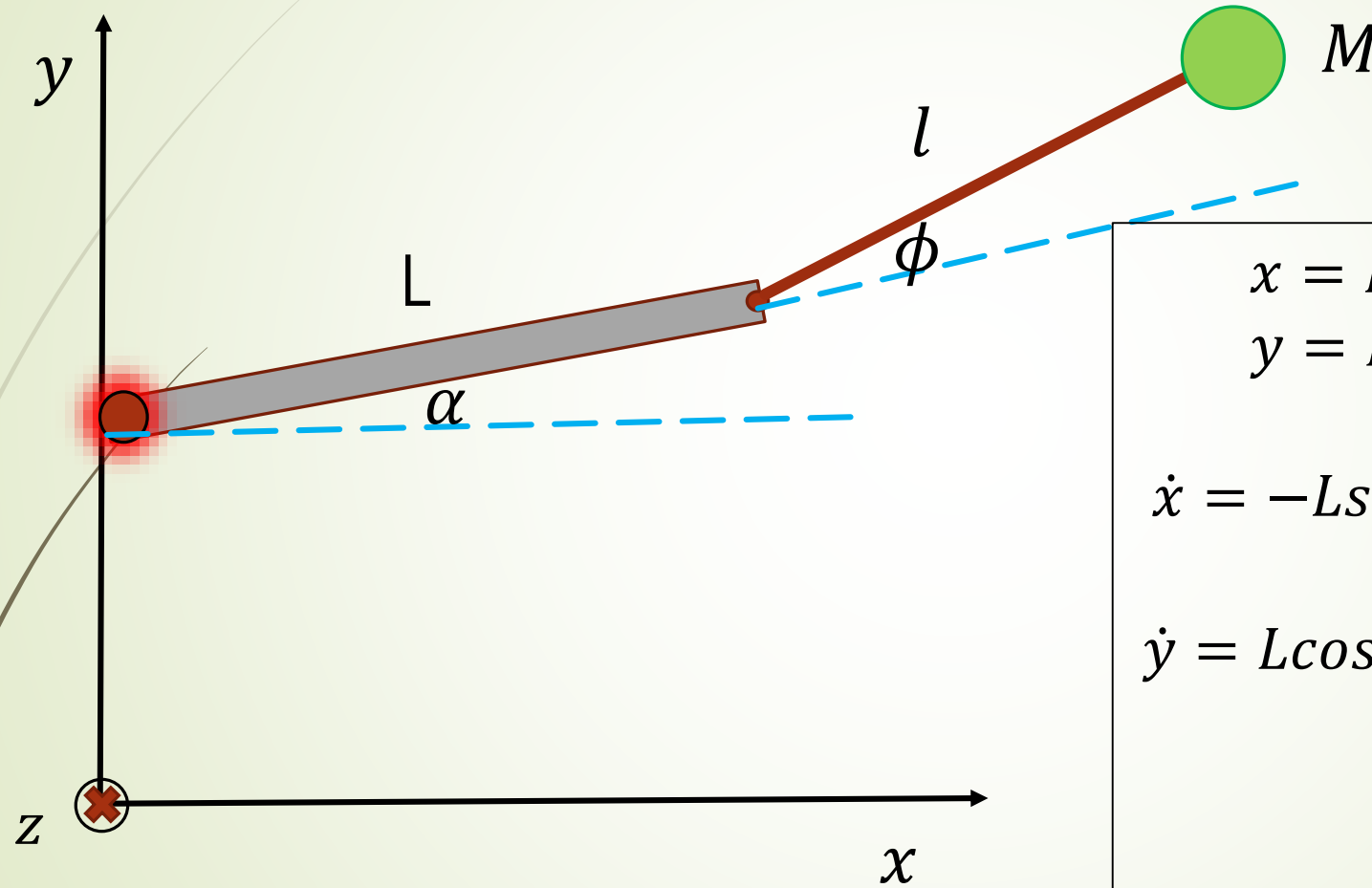
*Behaves like a simple
conical pendulum when
disturbed in the α direction*

Theory-Assumptions of the Model



- Ignoring the minimal deflection of the rod in the z -direction
- Assuming that l_2 remains constant (no contraction of the string)
- Assuming that the rod traces a circular motion α

Defining the generalized co-ordinates



$$x = L\cos(\alpha) + l\sin(\theta)\cos(\phi)$$

$$y = L\cos(\alpha) + l\sin(\theta)\sin(\phi)$$

$$z = -L\cos(\theta)$$

$$\dot{x} = -L\sin(\alpha)\dot{\alpha} + l(-\sin(\theta)\sin(\phi)\dot{\phi} + \cos(\theta)\cos(\phi)\dot{\theta})$$

$$\dot{y} = L\cos(\alpha)\dot{\alpha} + l(\sin(\theta)\cos(\phi)\dot{\phi} + \cos(\theta)\sin(\phi)\dot{\theta})$$

$$\dot{z} = l\sin(\theta)\dot{\theta}$$

Lagrangian of the System

Kinetic Energy:

$$\frac{1}{2} M (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) + \frac{1}{2} \left(\frac{1}{3} M L^2 \right) \dot{\alpha}^2$$

Velocity of the bob
Velocity of the rod

Potential Energy:

$$-MgL \cos(\theta) + \frac{1}{2} k \alpha^2$$

Potential Energy of bob
Elastic potential of the rod

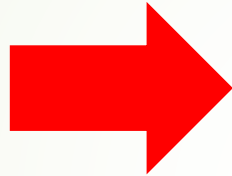
Lagrangian= K.E-P.E

Euler-Lagrange Equations

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = \frac{\partial L}{\partial \theta} \quad (1)$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\phi}} \right) = \frac{\partial L}{\partial \phi} \quad (2)$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\alpha}} \right) = \frac{\partial L}{\partial \alpha} \quad (3)$$



$$ML(\ddot{y} \cos(\alpha) - \ddot{x} \sin(\alpha)) = -k\alpha + \frac{1}{3}ML^2\ddot{\alpha}$$

$$\ddot{y} = \ddot{x} \tan(\phi)$$

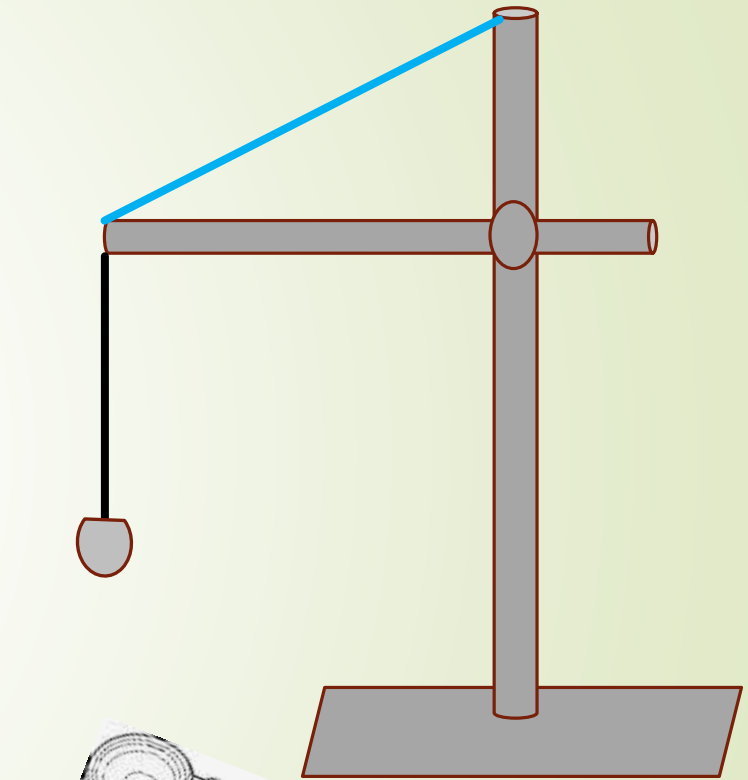
$$Ml\ddot{x} \cos(\phi) \cos(\theta) + Ml\ddot{y} \sin(\phi) \cos(\theta) + \ddot{z} \sin(\theta) = -Mgl \sin(\theta)$$

Can be written as 2nd order differential equations describing the acceleration in the (alpha), (theta) and (phi) directions

Governing Equations

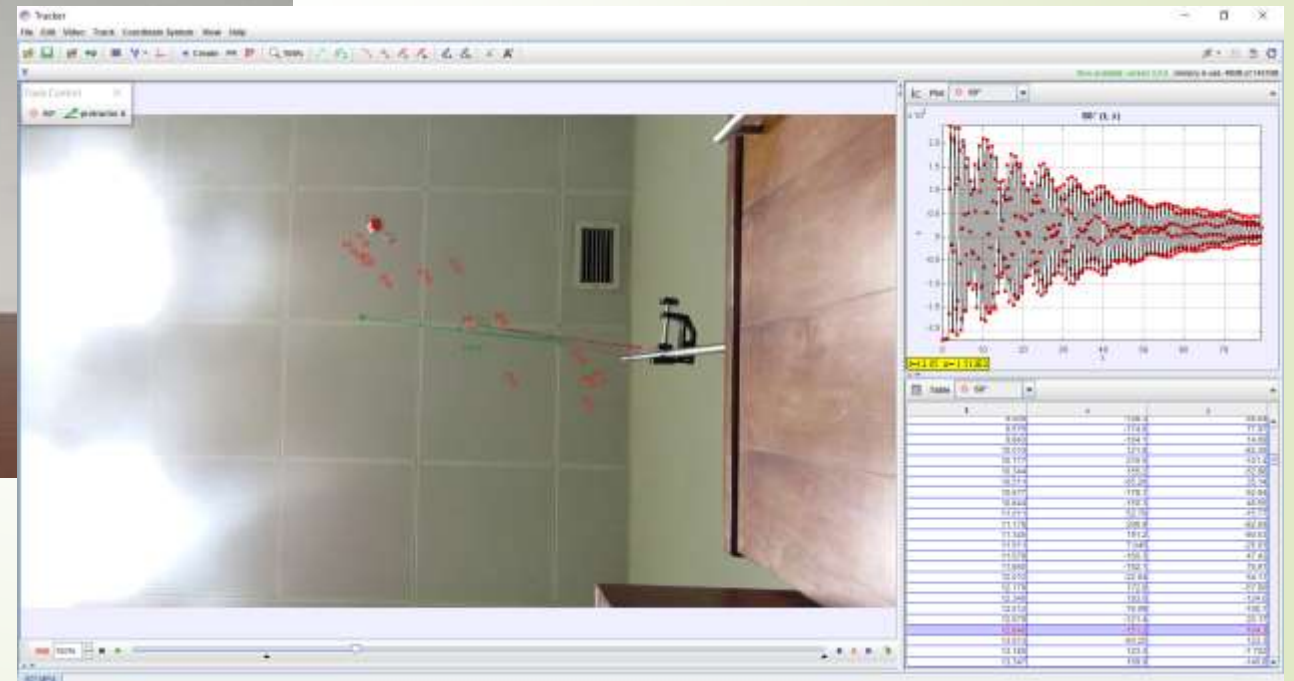
Experimental Set-up

- Rigid stand
- Strings
- Bobs of variable masses
- Camera
- Release mechanism for bob

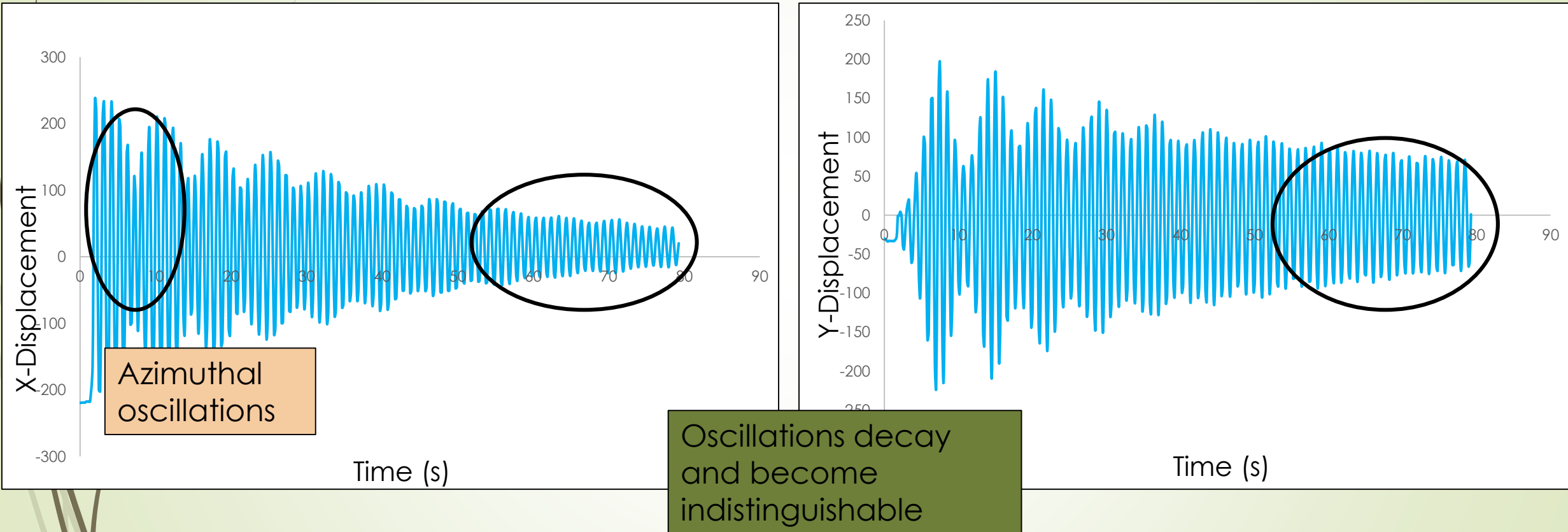


Camera

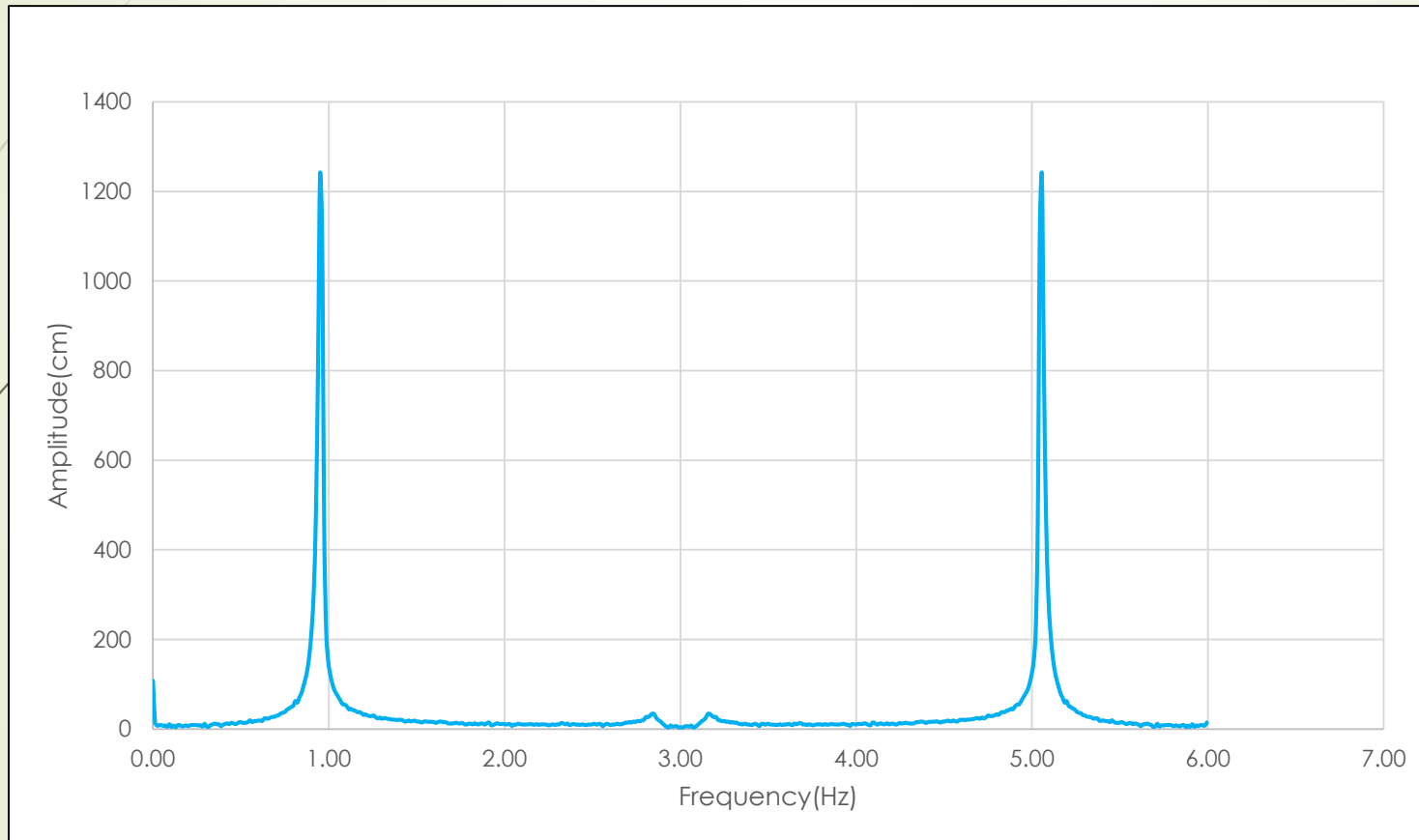
Experimental Set-up



Displacement-Time Graphs



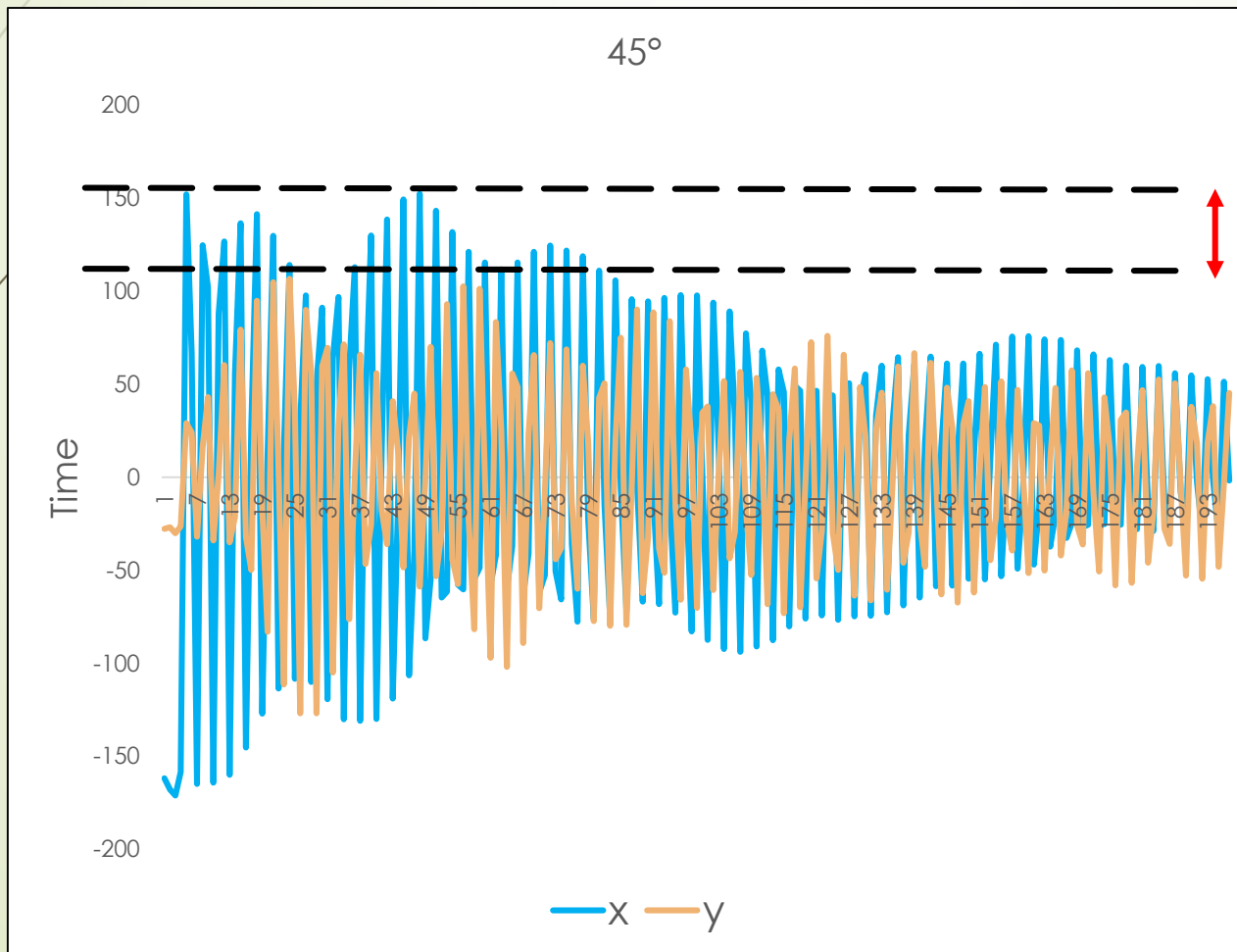
Fourier Analysis



Two bands corresponding to the
Azimuthal and Radial Oscillations

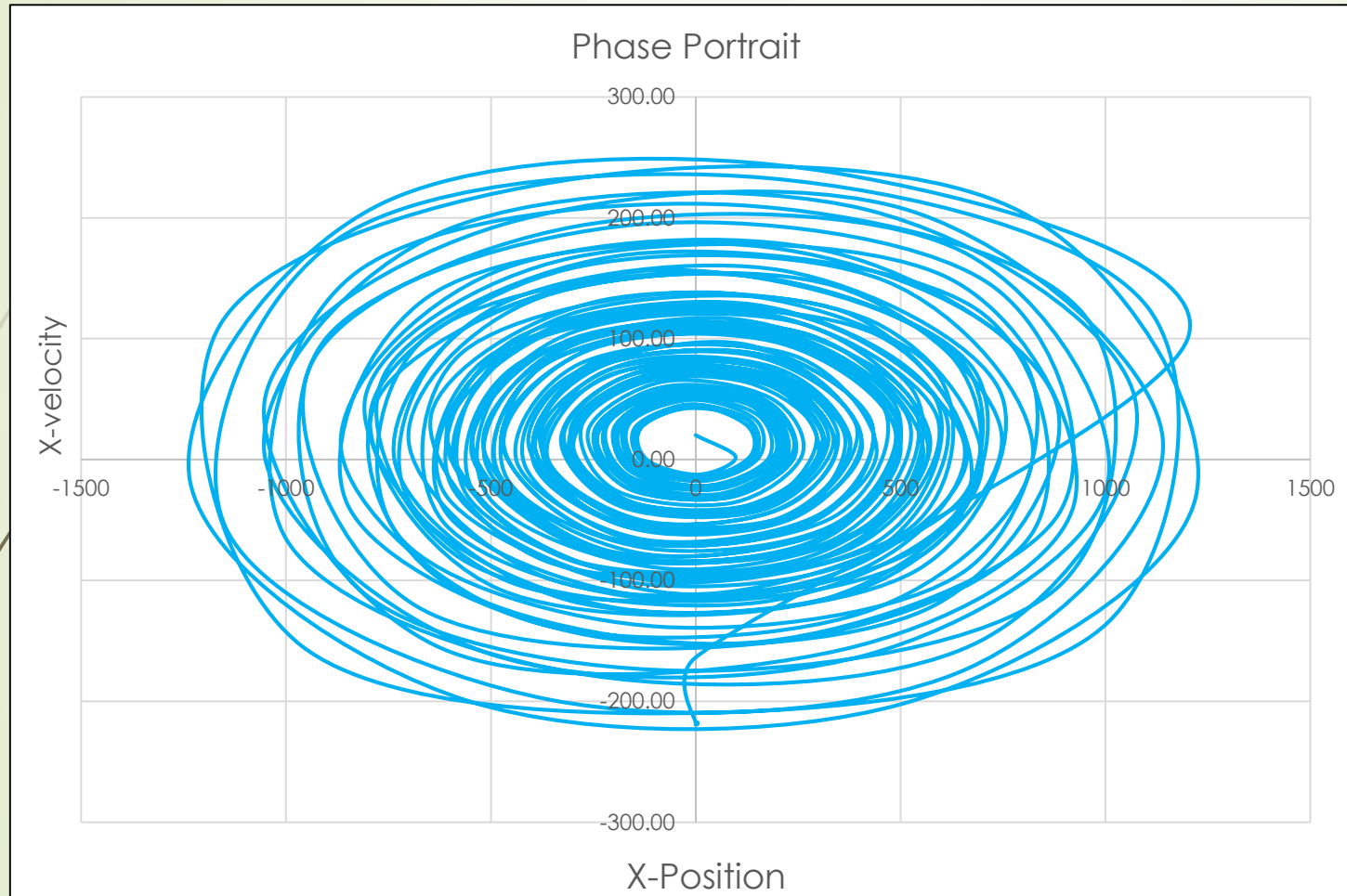
Comparing the Amplitudes

$$\text{Amplitude Ratio} = \frac{\text{Max amplitude in X direction}}{\text{Max amplitude in the Y direction}}$$



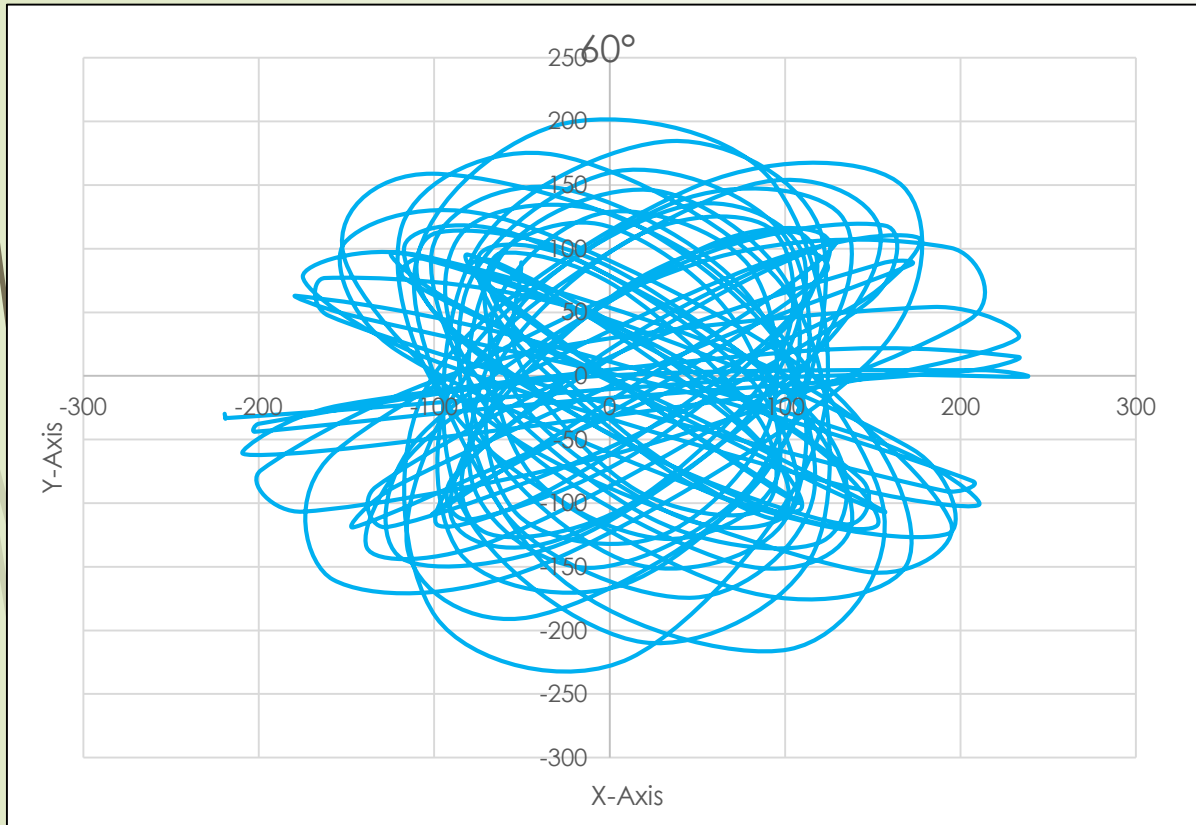
Obtained using the Max function in Excel

Phase Portrait

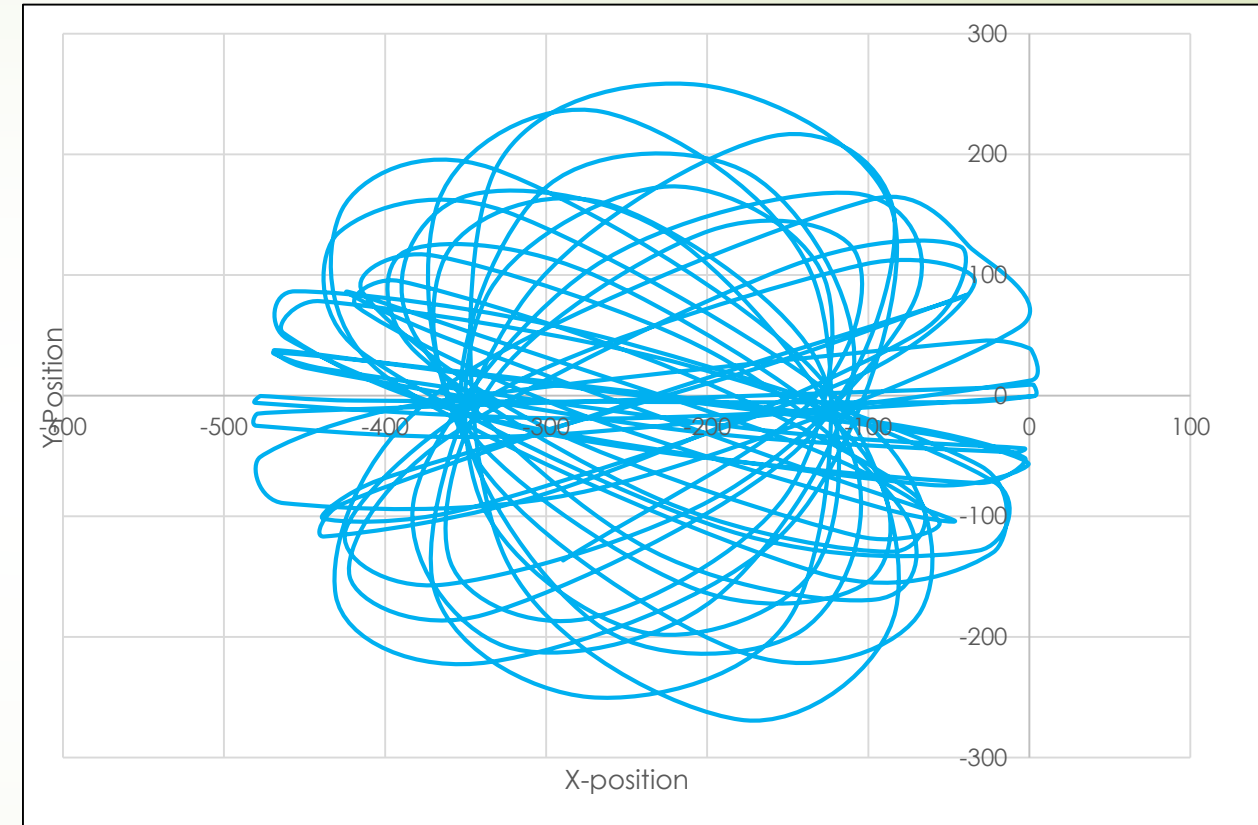


Shows non-chaotic, periodic motion for the general case

Experiment- release angle



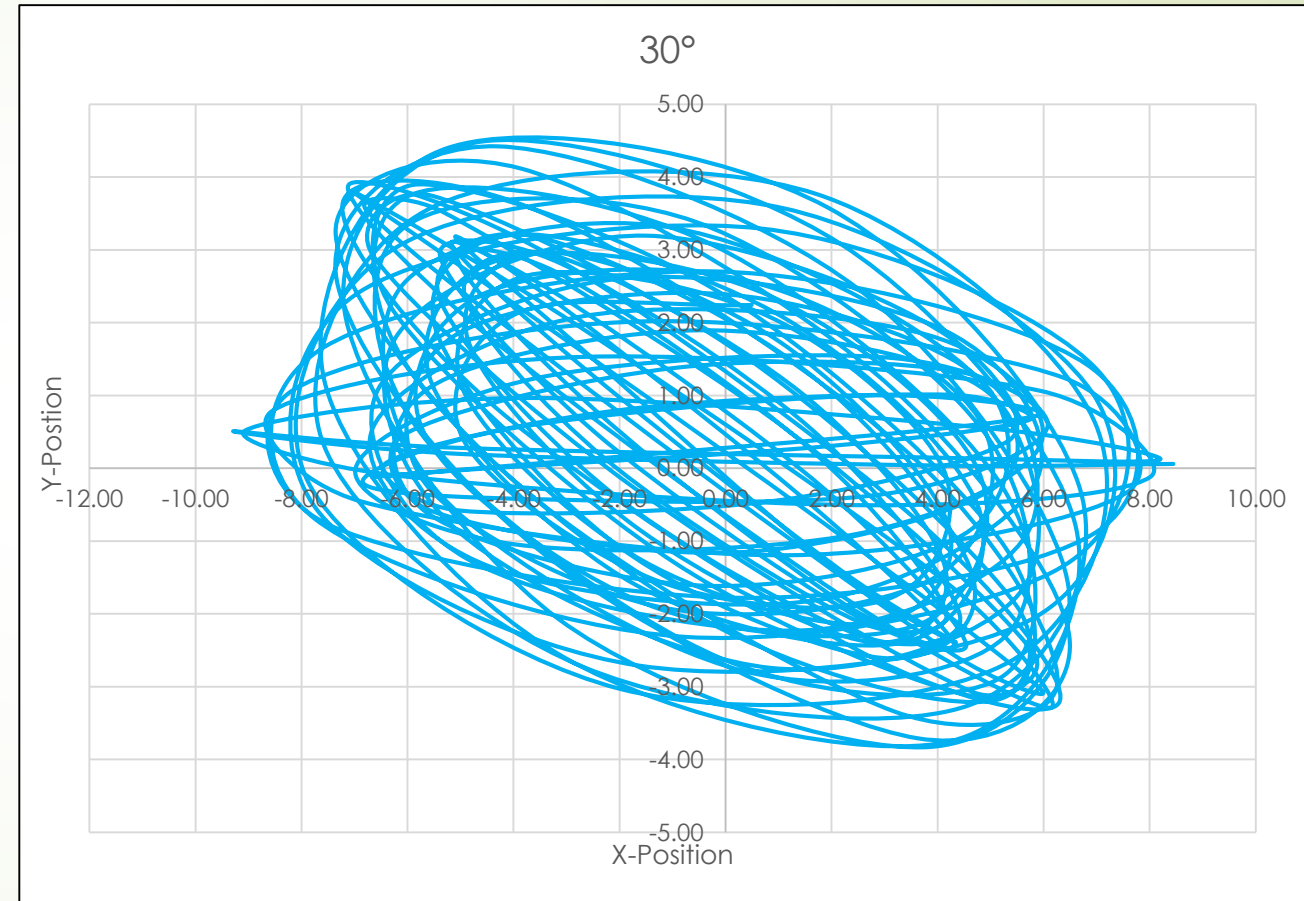
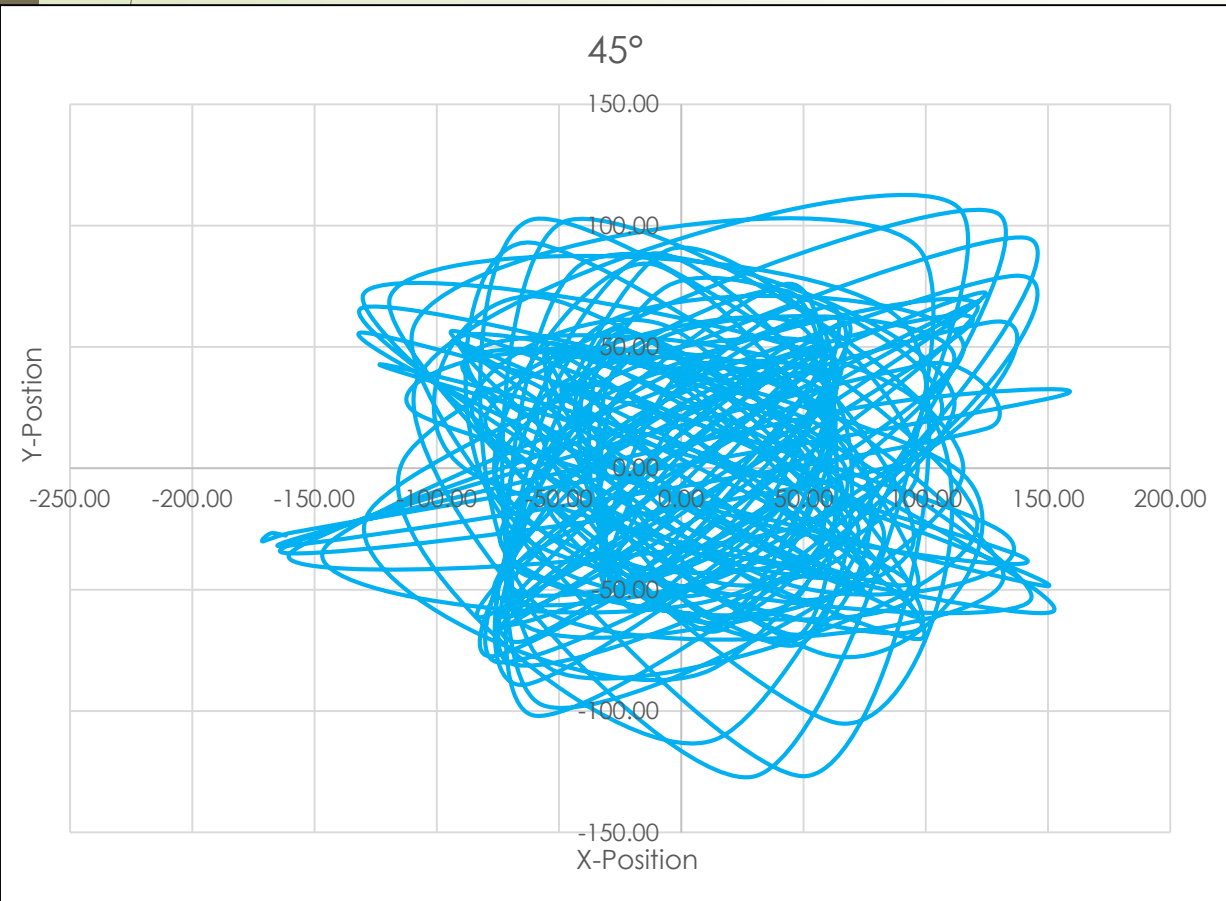
60°



90°

Bigger angles => higher transitions
rate

Initial Angles

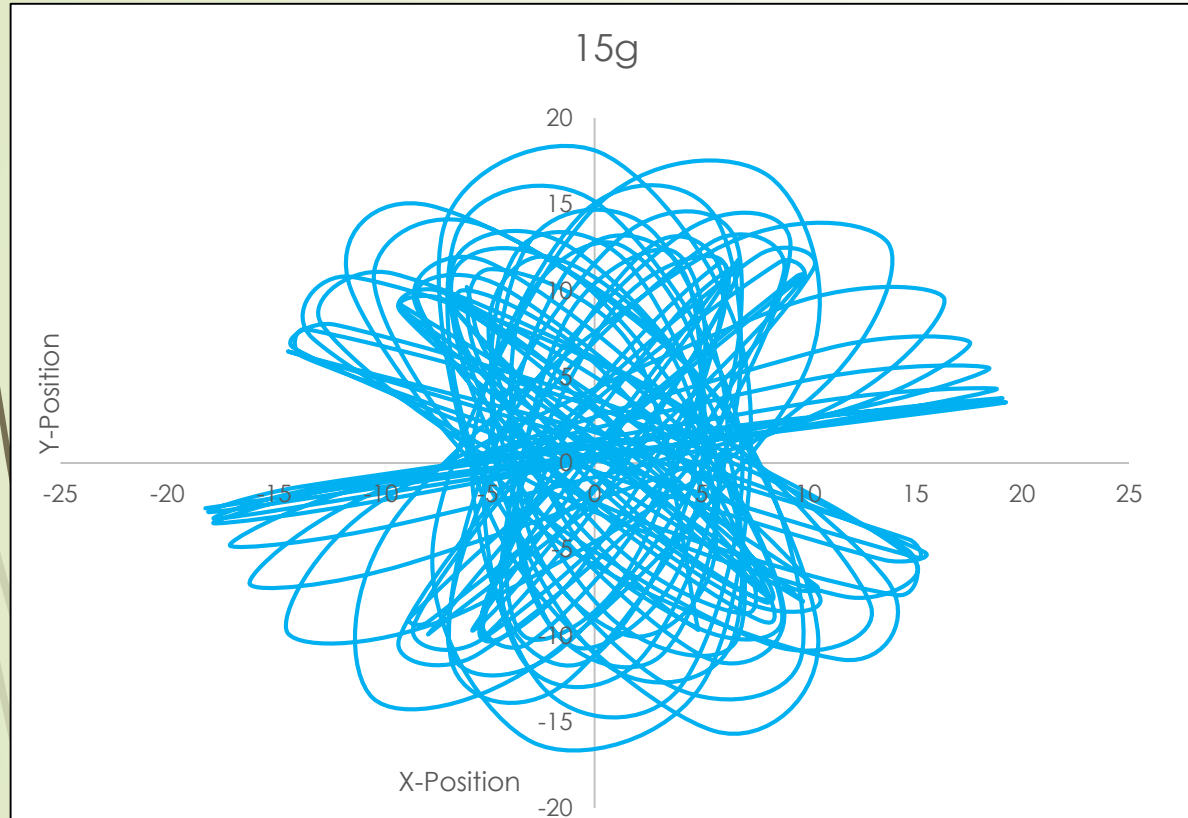


45°

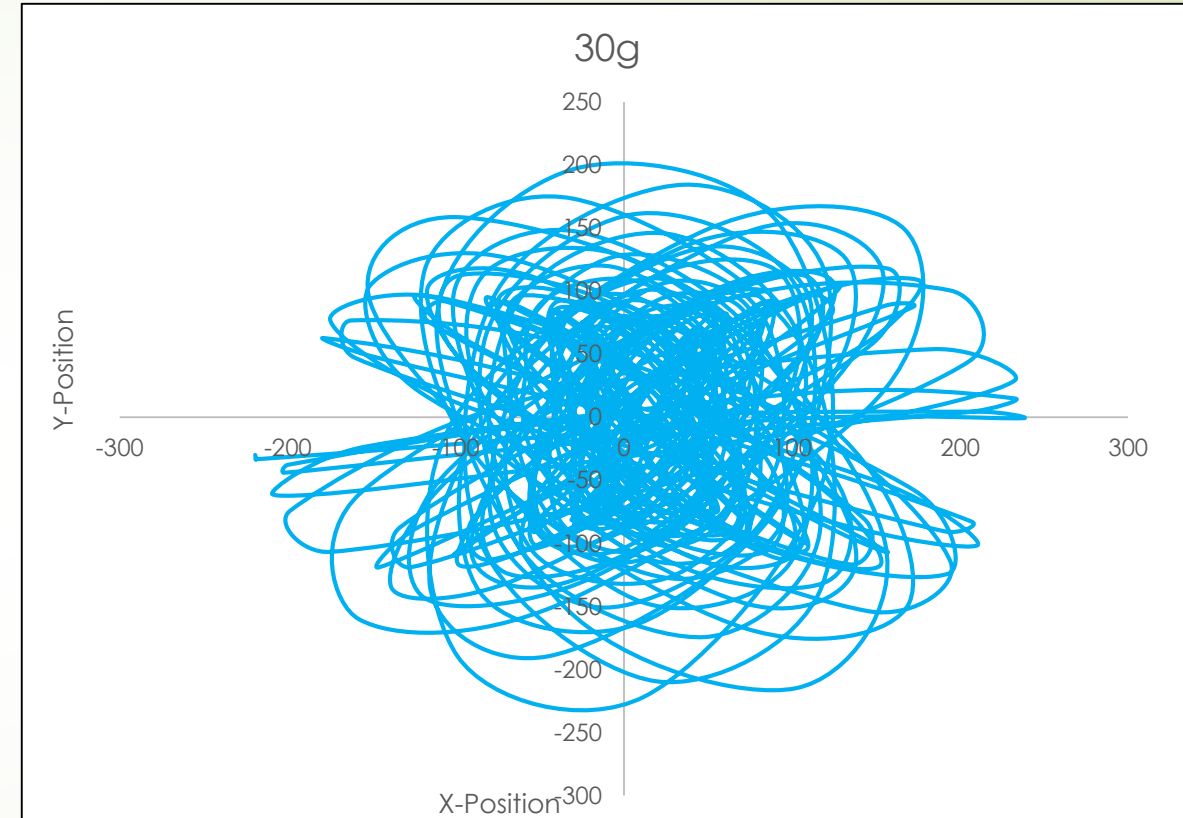
Smaller angles => Lower transitions rate

30°

Change in Mass (light)

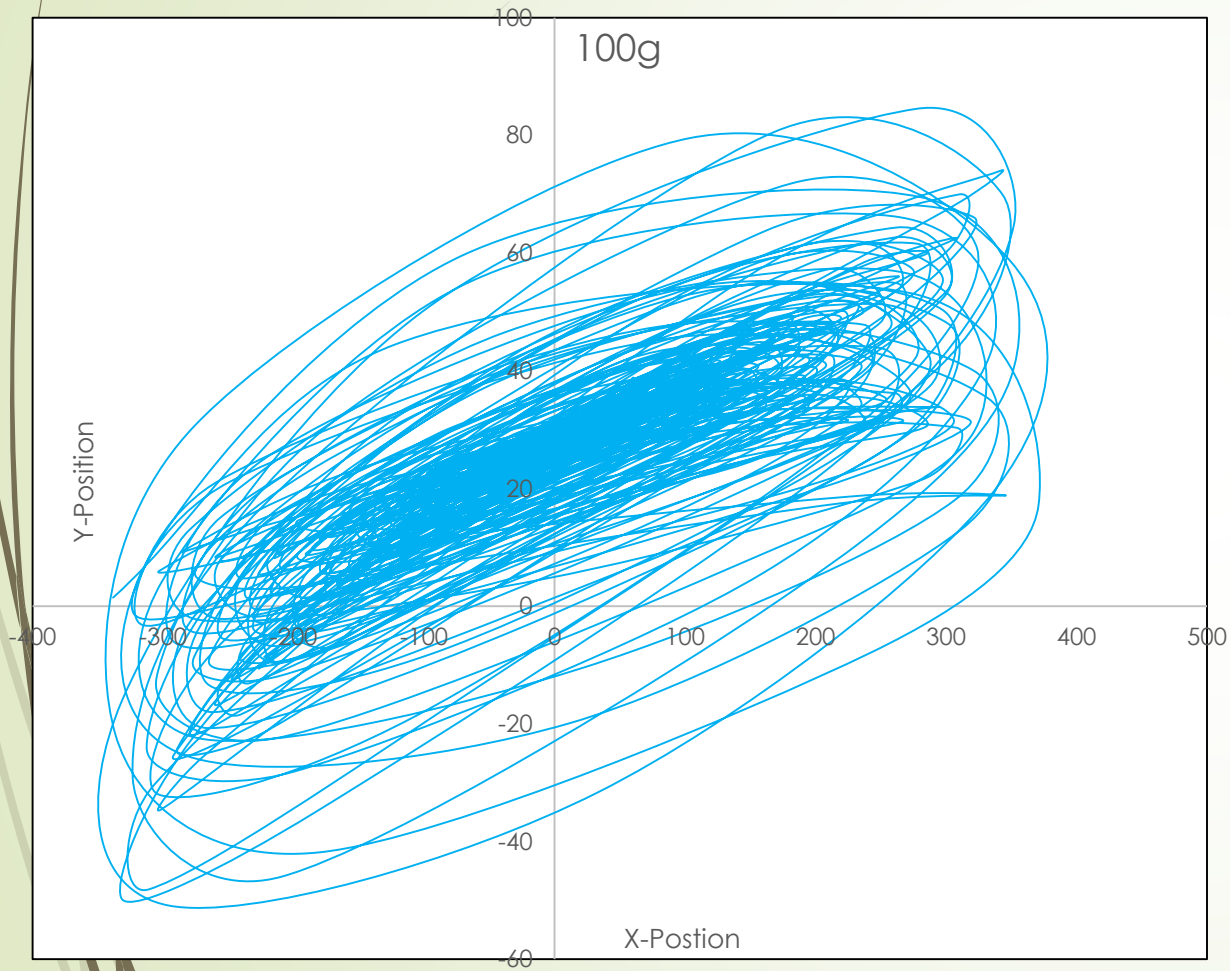


15g

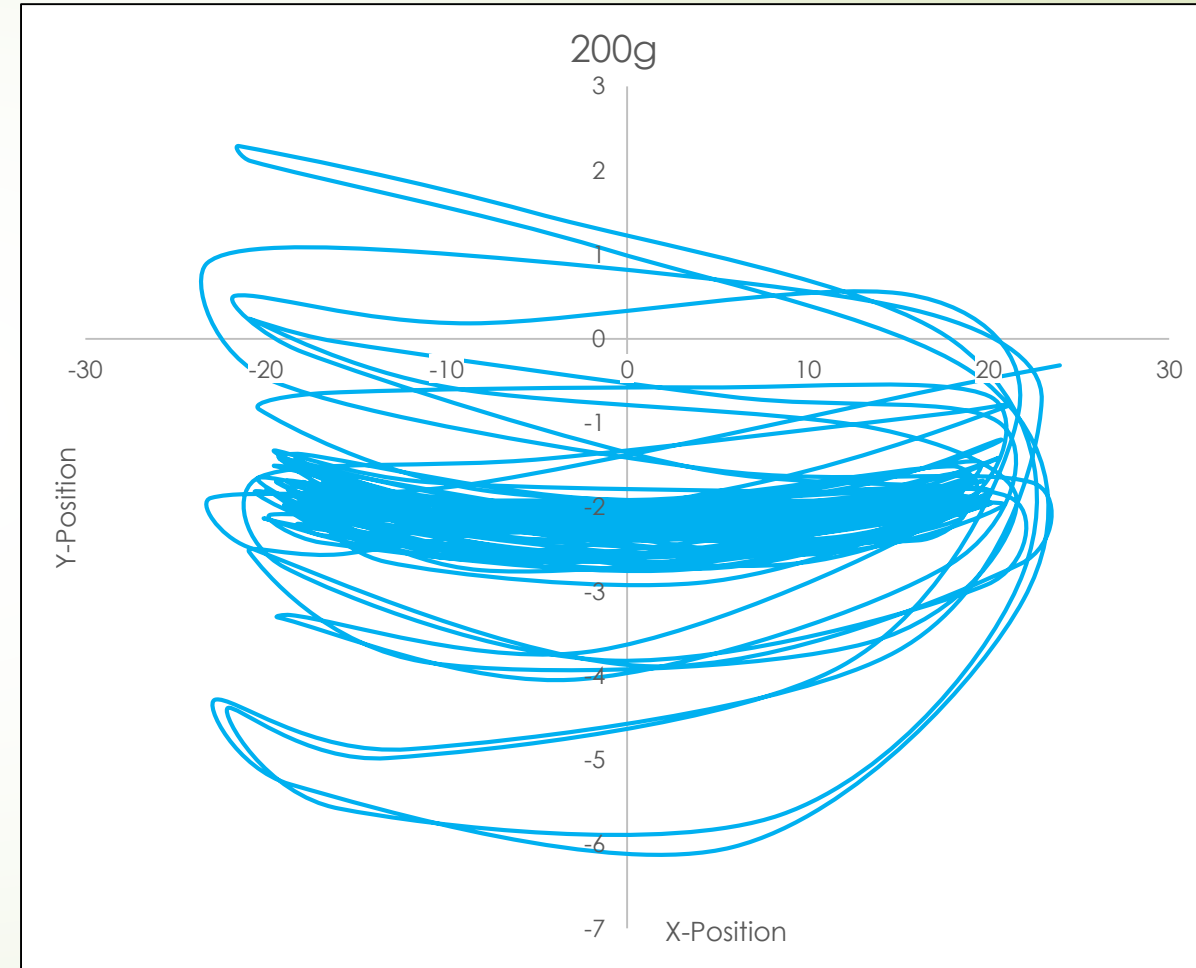


30g

Change in Mass (heavy)

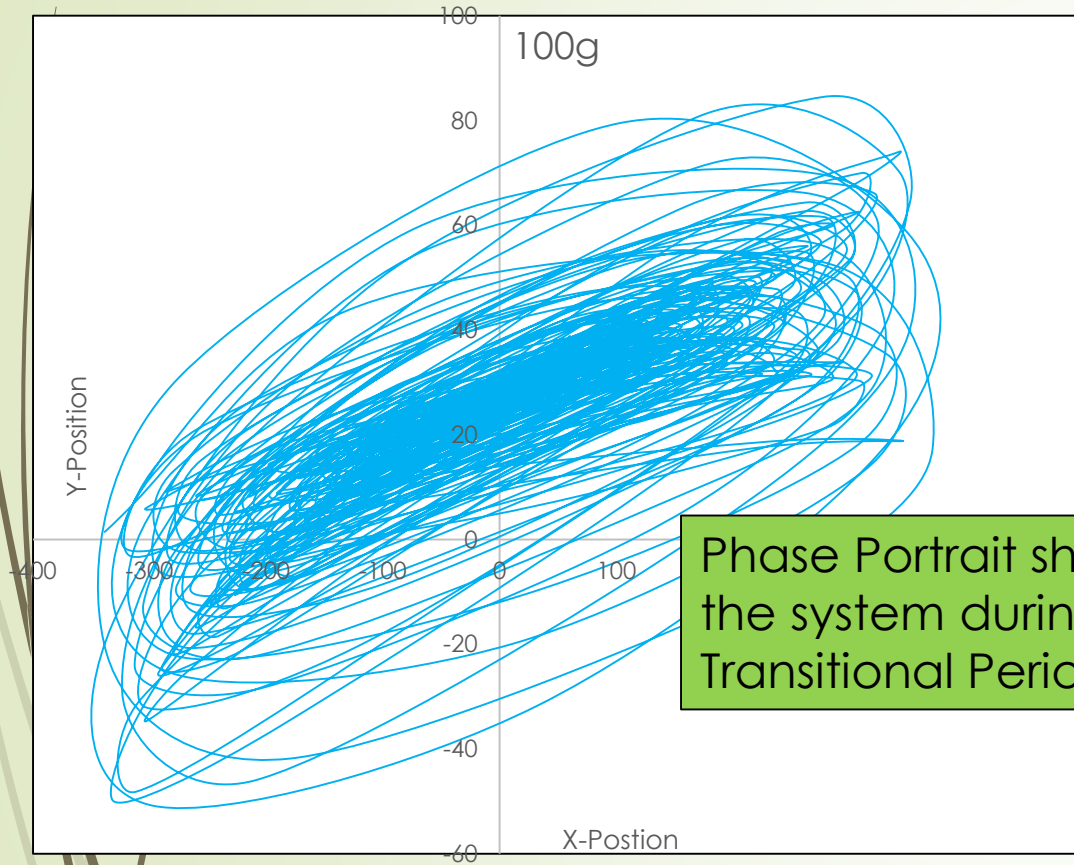


100g

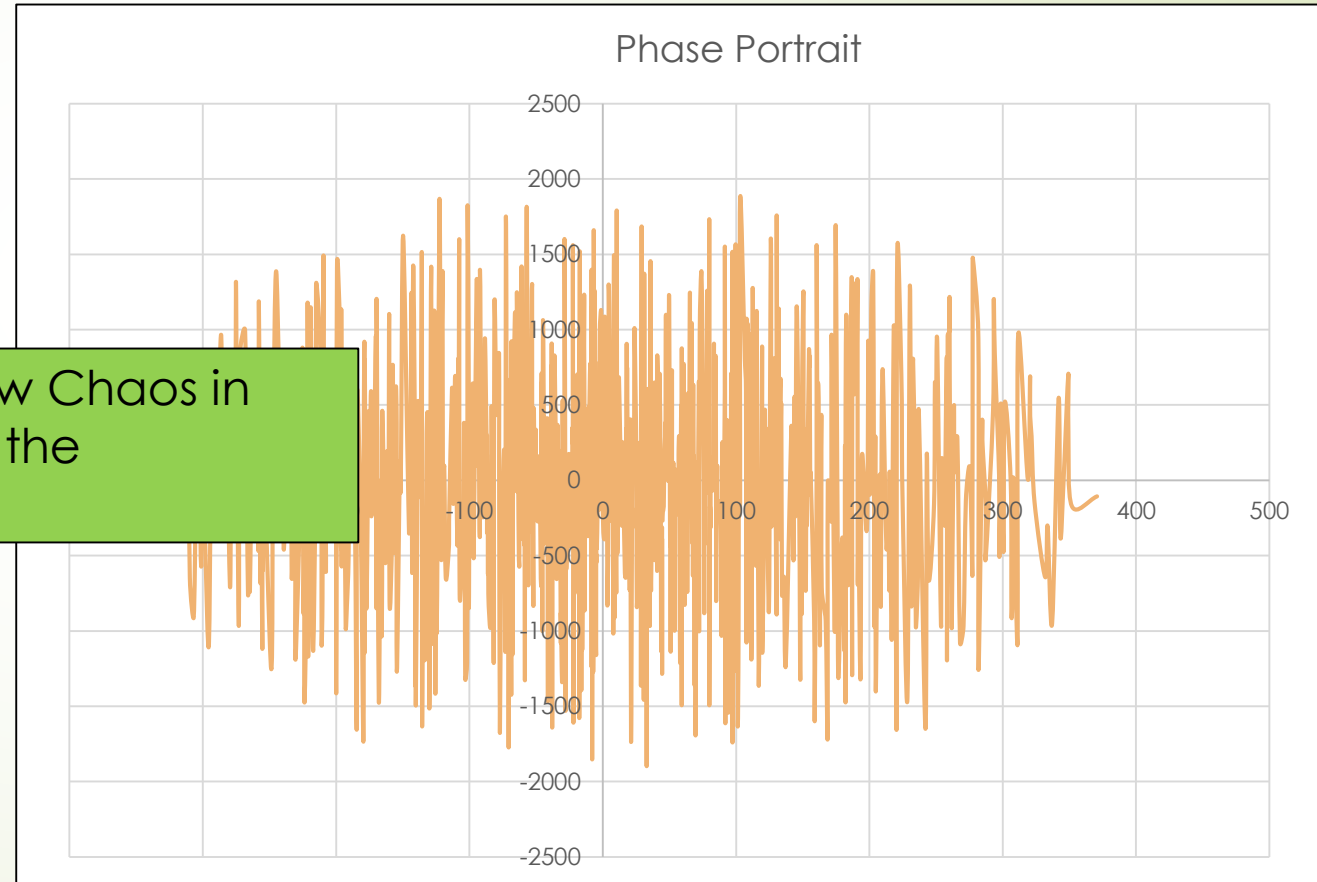


200g

Effect of Mass on the Elasticity of rod



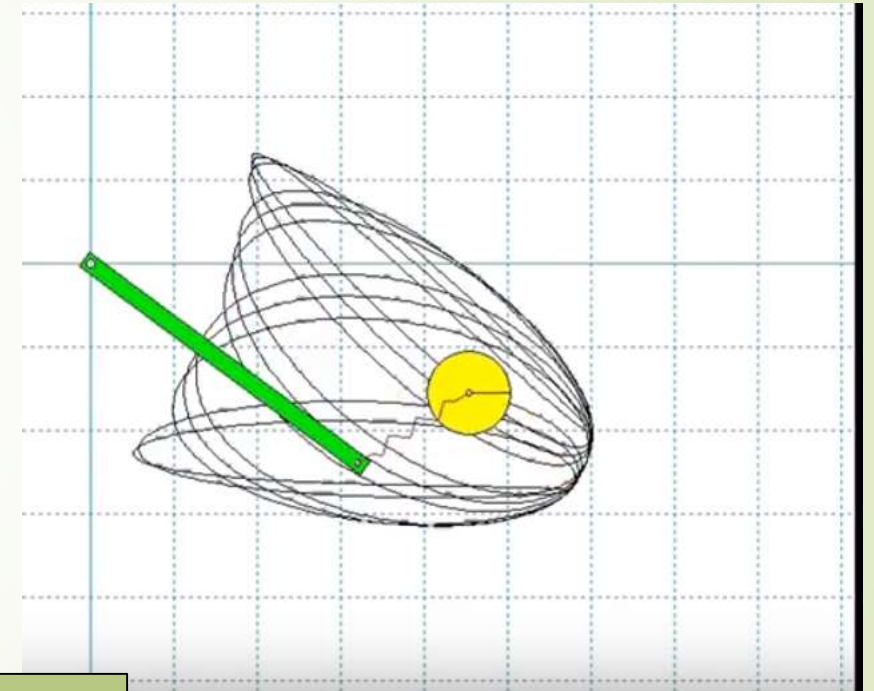
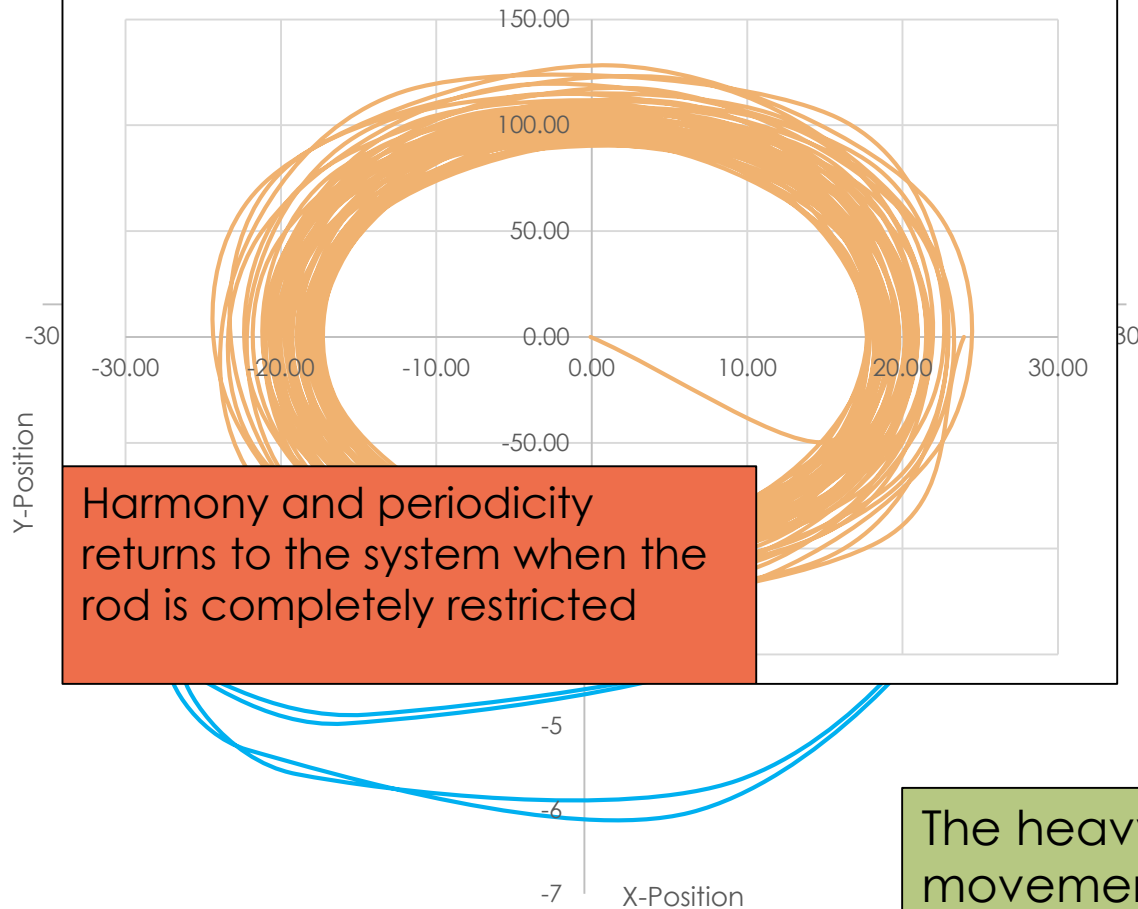
Phase Portrait show Chaos in the system during the Transitional Period



26

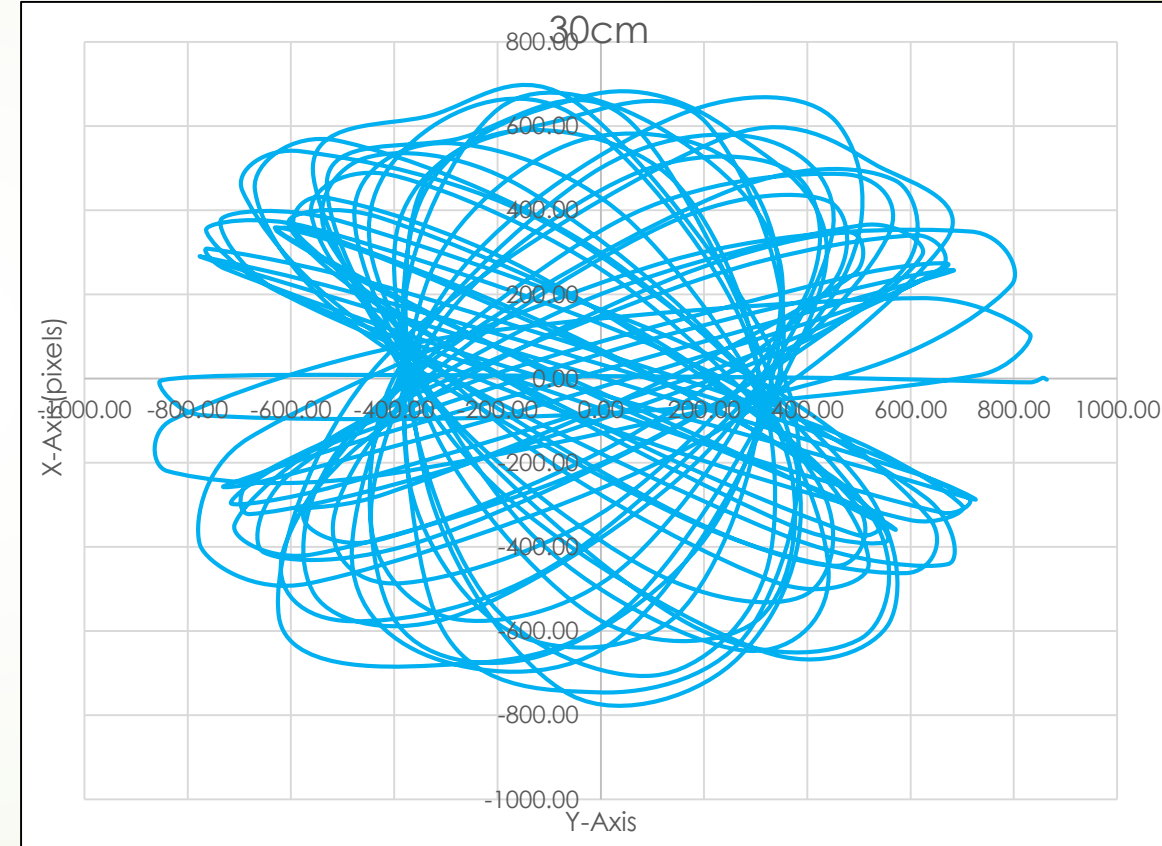
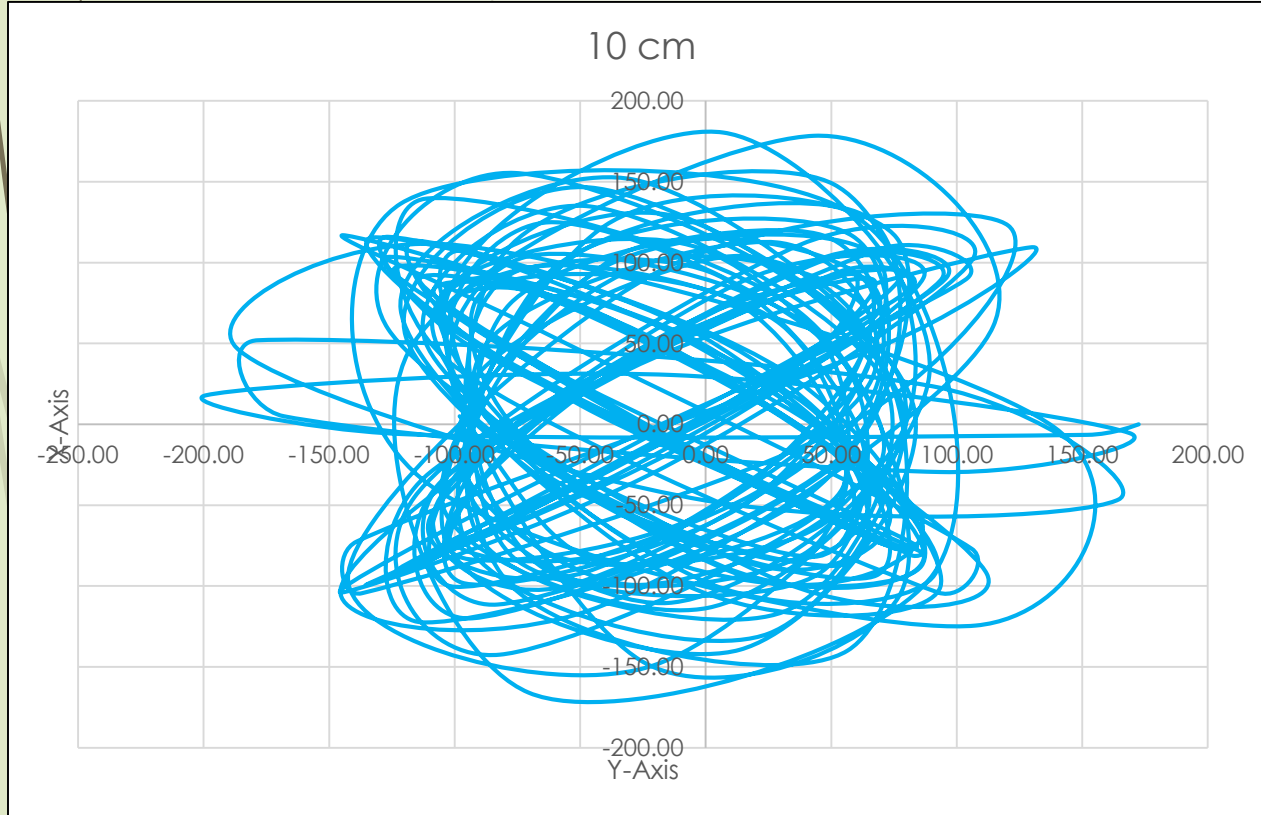
Effect of mass on the elasticity of rod

Phase Portrait



The heavy mass limits the movement of rod, leaving an almost rigid rod that can rotate (but not bend)

Length of the String

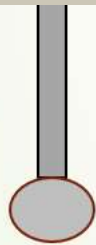


Conclusion-Preliminary

Observation 8 Investigation- Conditions of the Experiment

➤ Successfully described the Initial Conditions of the pendulum system

The pendulum
conical pen
oscillating fu



Rad **The pattern isn't observed when
disturbed in the phi direction**



Pendulum only
works when
disturbed in
the Radial
Direction

Conclusion-Theory

Theory- Lagrangian of the pendulum

y₃

Euler-Lagrange Equations

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\Theta}} \right) = \frac{\partial L}{\partial \Theta} \quad (1)$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\alpha}} \right) = \frac{\partial L}{\partial \alpha} \quad (2)$$

$$ML(\ddot{y} \cos(\alpha) - \ddot{x} \sin(\alpha)) = -k\alpha + \frac{1}{3}ML^2\ddot{\alpha}$$

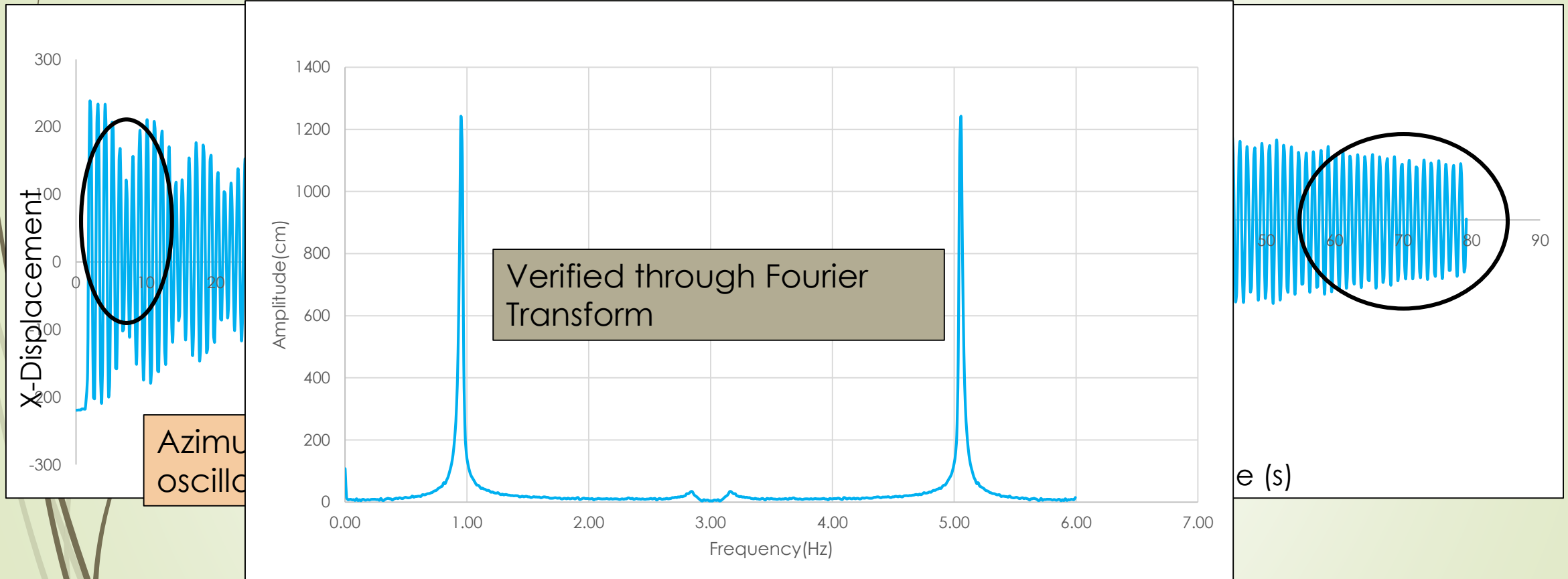
Successfully Modeled and defined
the System Mathematically

cos(θ)

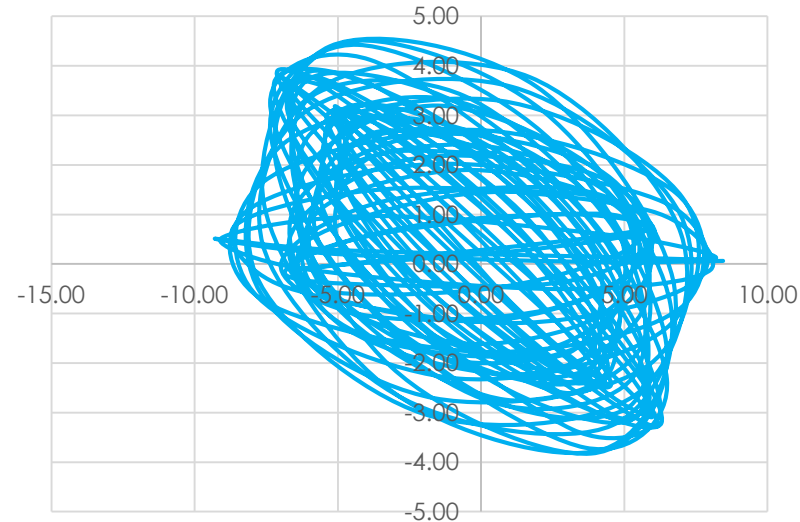
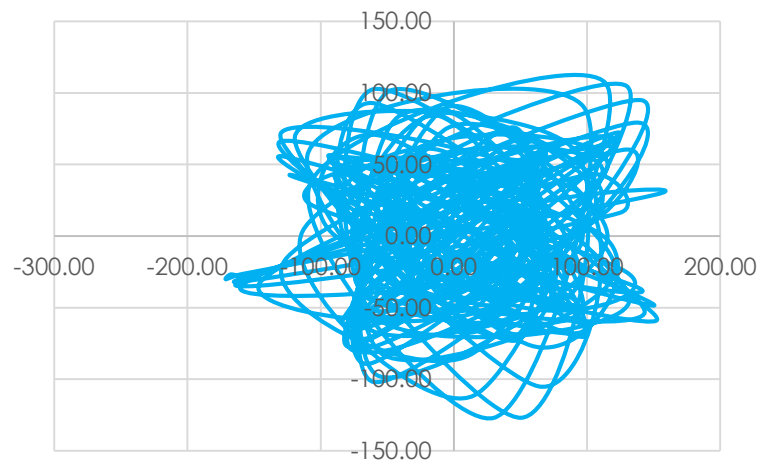
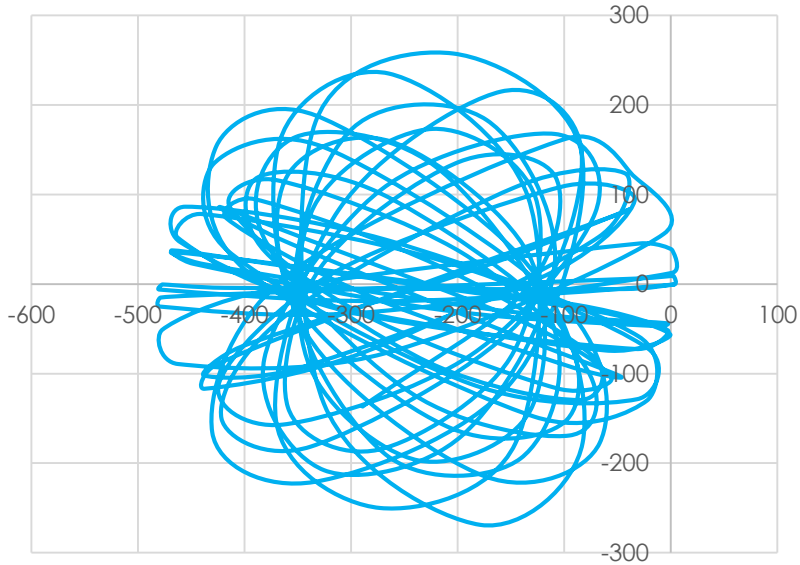
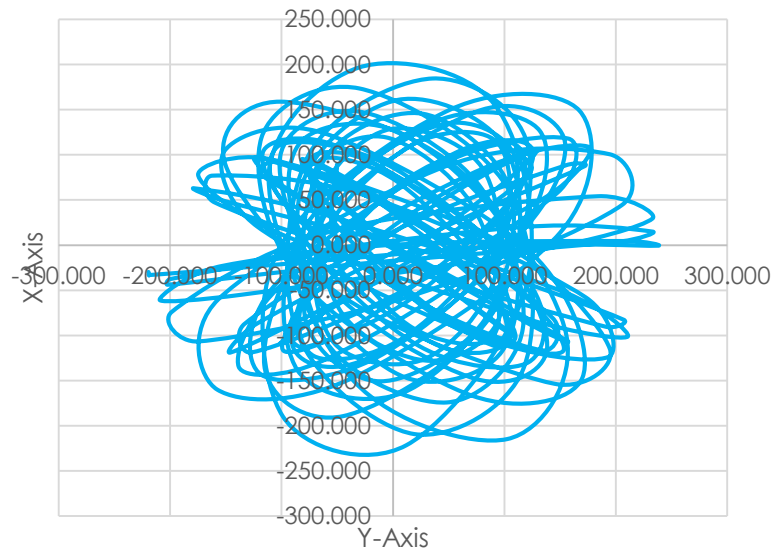
Governing Equations

The bob has 3
degrees of
Freedom which
complicate the
model

Conclusion-Experiments(Describing Frequencies)



Conclusion-Experiments



Investigated
and Explained
relations of the
system to 3
critical
parameters

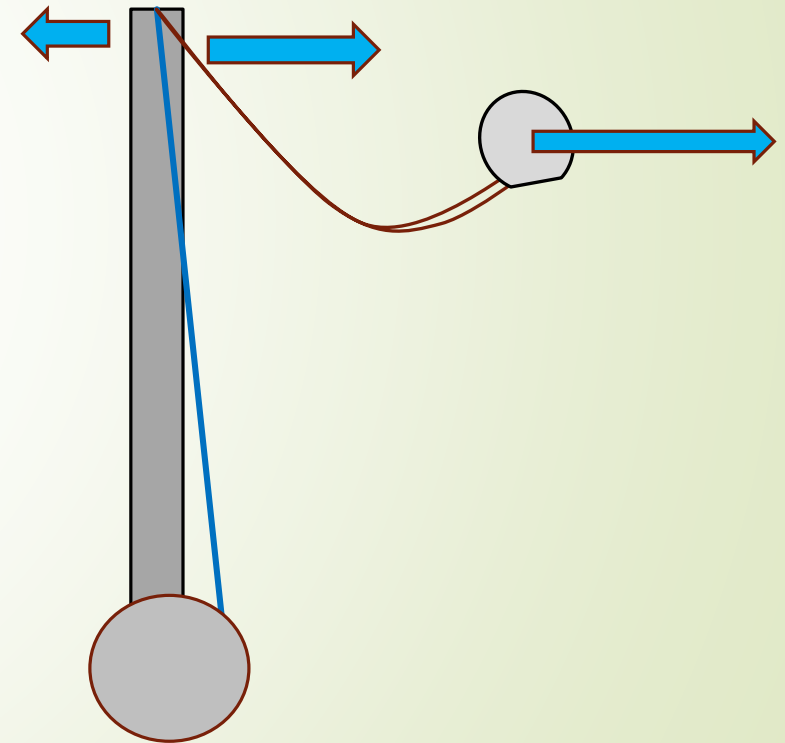
Thankyou!

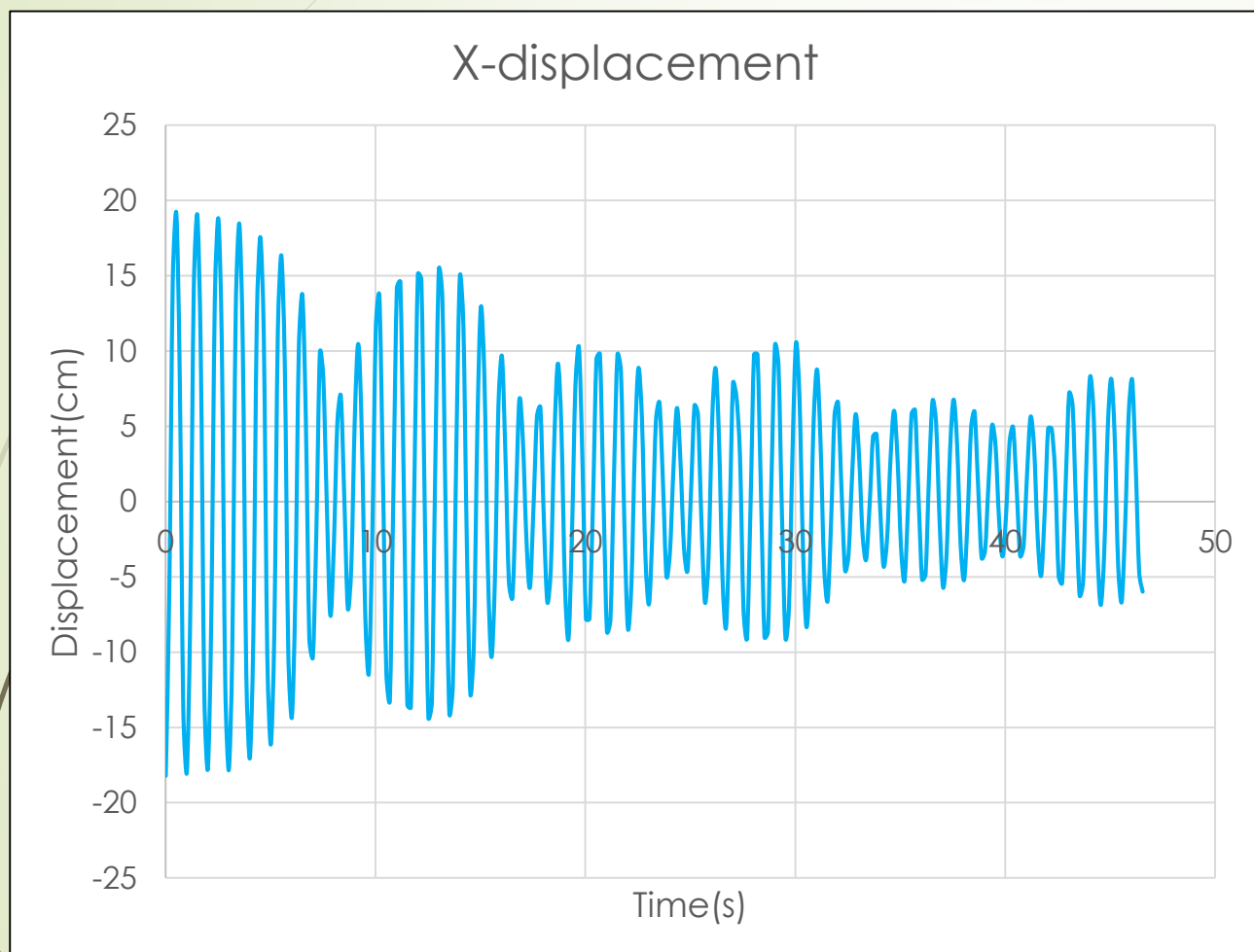
Relevant Parameters

- Length of the string
- Elasticity of the rod
- Weight of the bob
- Initial Angles

Theory- Conservation of energy

- ▶ The radial oscillations of the bob disturb the equilibrium of the elastic rod causing to move in the direction of the bob's motion.
- ▶ Since the rod is fixed on one end, it traces a circular path (azimuthal oscillations)
- ▶ The restoring force of the rod acts against the motion of the bob.
- ▶ This sets up a chain of energy transfer where the motion of the bob is transferred into the azimuthal perturbations of the rod





$\text{Mol. Weight} = 116.07 \text{ g/mol}$
 $\text{Mol. Weight} = 116.07$
 $\text{Mol. Weight} = 116.07$

$$M_{\text{eff}}^2 = \left(\frac{m_{\text{eff}}^2}{2\pi\alpha'} \right) = -\log \left[-\frac{1}{2} \frac{1}{\sqrt{1-\beta^2}} \right] = -\log \left[\frac{1}{2} \frac{1}{\sqrt{1-\beta^2}} \right]$$

$$\frac{1}{2} \left(\frac{1}{2} (2 \cos^2(\theta) - 1) \right) = \frac{1}{4} (2 \cos^2(\theta) - 1)$$

$$\frac{2}{\pi} = \frac{M_0 \left(\sin^2 \theta \cos^2 \theta - \sin^2 \theta \cos^2 \theta \right)}{\pi \left(\sin^2 \theta \cos^2 \theta + \sin^2 \theta \cos^2 \theta \right)}$$

$$= \frac{1}{n} \sum_{j=1}^n (\bar{x}_j - \bar{\bar{x}})^2 = \frac{1}{n} \sum_{j=1}^n (\bar{x}_j^2 - 2\bar{x}_j\bar{\bar{x}} + \bar{\bar{x}}^2) = \frac{1}{n} \left(\sum_{j=1}^n \bar{x}_j^2 - 2\bar{\bar{x}} \sum_{j=1}^n \bar{x}_j + n\bar{\bar{x}}^2 \right)$$

$$+ \frac{1}{2} \ln \left(\frac{1 + \sqrt{1 + 4x}}{1 - \sqrt{1 + 4x}} \right) + \frac{1}{2} \ln \left(\frac{1 + \sqrt{1 + 4x}}{1 - \sqrt{1 + 4x}} \right)$$

$$f(x) = \frac{1}{2} \ln(x^2 + 1) - \frac{1}{2} \ln(x^2 - 1)$$

$$\frac{f}{f} = \frac{f_{\text{max}}}{f_{\text{max}}} + \frac{f_{\text{min}}}{f_{\text{max}}} + \frac{f_{\text{avg}}}{f_{\text{max}}} + \frac{f_{\text{avg}}}{f_{\text{max}}}$$

$$\frac{1}{2} \cdot \ln(2x^2 + 3x^2) + \frac{1}{2} \cdot \ln\left(\frac{2x^2 + 3x^2}{2}\right)$$

$$N = m \sqrt{L \omega^2} + \frac{d}{dt} \sin \theta$$

$$T = T - V$$

$$E = \frac{1}{2}mv^2 + \left(\frac{p}{h}\right)^2 \lambda^2$$

$$\frac{1}{9}(\cos 2\theta)(\sin 4\theta)^2 - \frac{1}{9}(\sin 2\theta)(\cos 4\theta)^2$$

$$= (-\sin \theta) \times (\cos \theta) \frac{d\theta}{dx} = -(\frac{d}{dx} \sin \theta)$$

10-10-10

$$f(x) = \frac{1}{x^2} = x^{-2} \Rightarrow f'(x) = -2x^{-3} = -\frac{2}{x^3}$$

$$- \frac{1}{2} \frac{d}{dt} \left(\sin^2 \theta \right) \left(\frac{d\theta}{dt} \right) = \frac{1}{2} \frac{d}{dt} \left(\sin^2 \theta \right) \left(\frac{d\theta}{dt} \right)$$

$$+ \frac{1}{2} \ln \left(\frac{1 + \sqrt{1 - 4\alpha}}{1 - \sqrt{1 - 4\alpha}} \right)$$

$$\frac{d}{dt} \left(\int_{\Omega} u^2 dx \right) = -2 \int_{\Omega} u \Delta u dx$$

