

Q1. Eigenvalues of \hat{J}_z are $\hbar, 0, -\hbar$.

Eigenvalues of \hat{J}_z^2 are $\hbar^2, 0, \hbar^2$.

$$(a) \quad \hat{H} = -\frac{D}{\hbar} \begin{pmatrix} \hbar^2 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \hbar^2 \end{pmatrix}$$

$$= -D\hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

(b) If $D > 0$, ground state energy is $-D\hbar$.
If $D < 0$, ground state energy is 0 .

(c) Eigenstates of \hat{J}_z can be labeled as $|1, 1\rangle, |1, 0\rangle, |1, -1\rangle$, with $j = 1$, $m_j = 1, 0, -1$.

Let's find \hat{J}_+ in the $\{|1,1\rangle, |1,0\rangle, |1,-1\rangle\}$ basis.

$$\hat{J}_+ = \begin{pmatrix} \langle 1,1|\hat{J}_+|1,1\rangle & \langle 1,1|\hat{J}_+|1,0\rangle & \langle 1,1|\hat{J}_+|1,-1\rangle \\ \langle 1,0|\hat{J}_+|1,1\rangle & \langle 1,0|\hat{J}_+|1,0\rangle & \langle 1,0|\hat{J}_+|1,-1\rangle \\ \langle 1,-1|\hat{J}_+|1,1\rangle & \langle 1,-1|\hat{J}_+|1,0\rangle & \langle 1,-1|\hat{J}_+|1,-1\rangle \end{pmatrix}$$

$$\hat{J}_+ |1,0\rangle = \sqrt{2} \hbar |1,1\rangle$$

$$\begin{aligned} \hat{J}_+ |1,-1\rangle &= \sqrt{2 - (-1)(0)} \hbar |1,0\rangle \\ &= \sqrt{2} \hbar |1,0\rangle. \end{aligned}$$

$$\hat{J}_+ |1,-1\rangle = \sqrt{2} \hbar |1,0\rangle$$

$$\therefore \hat{J}_+ = \hbar \begin{pmatrix} 0 & \sqrt{2} & 0 \\ 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix}$$

$$\hat{J}_- = \hat{J}_+^\dagger = \hbar \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 \\ 0 & \sqrt{2} & 0 \end{pmatrix}$$

$$(d) \quad \hat{J}_+^2 = \hbar^2 \begin{pmatrix} 0 & \sqrt{2} & 0 \\ 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & \sqrt{2} & 0 \\ 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix} \quad \text{Sp}$$

$$= \hbar^2 \begin{pmatrix} 0 & 0 & 2 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\hat{J}_-^2 = \hbar^2 \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 \\ 0 & \sqrt{2} & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & 0 \\ \sqrt{2} & 0 & 0 \\ 0 & \sqrt{2} & 0 \end{pmatrix}$$

$$= \hbar^2 \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 2 & 0 & 0 \end{pmatrix}$$

$$\hat{J}_+^2 + \hat{J}_-^2 = 2\hbar^2 \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}$$

$$\hat{H}_2 = -D\hbar \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} + 2\hbar A \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix}$$

$$= \hbar \begin{pmatrix} -D & 0 & 2A \\ 0 & 0 & 0 \\ 2A & 0 & -D \end{pmatrix}$$

Eigenstates of \hat{H}_2 :

One of them is clearly $|1, 0\rangle$.

Pick up the block $\frac{1}{\hbar} \begin{pmatrix} -D & 2A \\ 2A & -D \end{pmatrix}$.

This looks like $-D \frac{\hat{1}}{\hbar} + \frac{2A}{\hbar} \hat{\sigma}_x$.

We know eigenstates of $\hat{\sigma}_x$ are $\frac{|a\rangle \pm |b\rangle}{\sqrt{2}}$. Let's

see if $\frac{|1, 1\rangle + |1, -1\rangle}{\sqrt{2}} = |+\rangle$ and $\frac{|1, 1\rangle - |1, -1\rangle}{\sqrt{2}} = |-\rangle$

qualify as eigenstates of this block.

$$\text{Note } \frac{|1, 1\rangle + |1, -1\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\frac{|1, 1\rangle - |1, -1\rangle}{\sqrt{2}} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

$$\frac{1}{\hbar} \begin{pmatrix} -D & 2A \\ 2A & -D \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} -D+2A \\ -D+2A \end{pmatrix} \\ = \frac{1}{\sqrt{2}} (-D+2A) \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

Hurray! $|1, 1\rangle + |1, -1\rangle$ is one eigenstate. The other one $\frac{1}{\sqrt{2}} (|1, 1\rangle - |1, -1\rangle)$ (must be).

Check:

$$\frac{1}{\hbar} \begin{pmatrix} -D & 2A \\ 2A & -D \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} -D-2A \\ 2A-D \end{pmatrix} \\ = \frac{1}{\sqrt{2}} (-D-2A) \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

Now initial state is

$$|1, -1\rangle = \frac{1}{\sqrt{2}} (|+\rangle - |-\rangle).$$

expressing as eigenstates of \hat{H}_2 .

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$$|\psi(t)\rangle = e^{-i\frac{\hat{H}t}{\hbar}} |1, -1\rangle$$

$$= e^{-i\frac{\hat{H}t}{\hbar}} \frac{1}{\sqrt{2}} (|+\rangle - |- \rangle)$$

$$= \frac{e^{-i(-D+2A)t} |+\rangle - e^{-i(-D-2A)t} |- \rangle}{\sqrt{2}}$$

$$= \frac{e^{iDt} e^{-i2At} |+\rangle - e^{iDt} e^{+i2At} |- \rangle}{\sqrt{2}}$$

$$= \frac{e^{-i2At} |+\rangle - e^{+i2At} |- \rangle}{\sqrt{2}}$$

$$= \frac{e^{-i2At}}{\sqrt{2}} \left(|1, 1\rangle + |1, -1\rangle \right)$$

$$\neq \frac{e^{+i2At}}{\sqrt{2}} \left(|1, 1\rangle - |1, -1\rangle \right)$$

~~$$= \frac{1}{\sqrt{2}} \left(2 \cos 2At \right) |1, 1\rangle$$~~

(ignore
global
phase)

Put things back
- perspective

$$= \frac{1}{2} \begin{pmatrix} e^{-i2At} & -e^{+i2At} \\ e^{-i2At} & +e^{+i2At} \end{pmatrix} \begin{pmatrix} 1, 1 \end{pmatrix} + \frac{1}{2} \begin{pmatrix} e^{-i2At} & -e^{+i2At} \\ e^{-i2At} & +e^{+i2At} \end{pmatrix} \begin{pmatrix} 1, -1 \end{pmatrix}$$

$$= -\frac{1}{2} \cdot 2i \sin(2At) \begin{pmatrix} 1, 1 \end{pmatrix} + \frac{1}{2} \cdot 2 \cos(2At) \begin{pmatrix} 1, -1 \end{pmatrix}$$

$$= -i \sin(2At) \begin{pmatrix} 1, 1 \end{pmatrix} + \cos(2At) \begin{pmatrix} 1, -1 \end{pmatrix}$$

Q2 The state $|a\rangle$ is an eigenstate of \hat{S}_x with eigenvalue $\frac{3}{2}\hbar$.
 Let's write this state in vector form.

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$$\hat{S}_x |a\rangle = \frac{3}{2}\hbar |a\rangle$$

$$\frac{1}{\hbar} \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ \sqrt{3} & 0 & 2 & 0 \\ 0 & 2 & 0 & \sqrt{3} \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{pmatrix} = \frac{3}{2} \begin{pmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{pmatrix}$$

$$\sqrt{3} c_2 = \frac{3}{2} c_1 \Rightarrow$$

$$c_2 = \frac{\sqrt{3}}{2} c_1$$

$$\sqrt{3} c_1 + 2 c_3 = \frac{3}{2} c_2$$

$$2 c_2 + \sqrt{3} c_4 = \frac{3}{2} c_3$$

$$\sqrt{3} c_3 = \frac{3}{2} c_4 \Rightarrow$$

$$c_3 = \frac{\sqrt{3}}{2} c_4$$

$$\Rightarrow \sqrt{3} \cdot \cancel{\sqrt{3}} c_1 + 2 c_3 = \frac{3}{2} \frac{\sqrt{3}}{2} c_1$$

$$2 C_3 = \left(3 \frac{\sqrt{3}}{4} - \sqrt{3} \right) C_1$$

$$C_3 = \left(\frac{3\sqrt{3}}{8} - \frac{\sqrt{3}}{2} \right) C_1$$

$$C_4 = \frac{2}{\sqrt{3}} \left(3 \frac{\sqrt{3}}{8} - \frac{\sqrt{3}}{2} \right) C_1$$

$$C_4 = \left(\frac{3}{4} - \sqrt{3} \right) C_1$$

$$|a\rangle = N \begin{pmatrix} 1 \\ \frac{\sqrt{3}}{2} \\ \frac{3\sqrt{3}}{8} - \frac{\sqrt{3}}{2} \\ \frac{3}{4} - \sqrt{3} \end{pmatrix}$$

N is a normalization constt.

✓ x

$$|\langle a|a \rangle|^2 = 1$$

$$= N^2 \left(1 + \frac{3}{4} + \left(\frac{3\sqrt{3}}{8} - \frac{\sqrt{3}}{2} \right)^2 + \left(\frac{3}{4} - \sqrt{3} \right)^2 \right)$$

$$= N^2 \left(1 + \frac{3}{4} + \frac{27}{64} + \frac{3}{4} - \frac{9}{8} + \frac{9}{16} + 3 - \frac{3\sqrt{3}}{2} \right)$$

$$= N^2 \left(\frac{64 + 48 + 27 + 48 - 72 + 36 + 192 - 96\sqrt{3}}{64} \right)$$

$$= N^2 \left(\frac{115}{64} + \left(\frac{3}{4} - \sqrt{3} \right)^2 \right)$$

$$= N^2 \left(\frac{343 - 96\sqrt{3}}{64} \right) = 1$$

$$\Rightarrow N^2 = \frac{64}{343 - 96\sqrt{3}}$$

$$N = \sqrt{\frac{64}{343 - 96\sqrt{3}}} = 0.6018$$

If the input is an equal superposition of $\mathbb{S}_0 \times$
all eigenstates of \underline{S}_x , $\frac{1}{4}$ of them will
 come out of the top channel.

If input state is $|\frac{3}{2}, \frac{3}{2}\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}$, then the

Fraction of states coming out will be

$$\left| N \begin{pmatrix} 1 & \frac{\sqrt{3}}{2} & \frac{3\sqrt{3}-\sqrt{3}}{8} & \frac{3-\sqrt{3}}{4} \end{pmatrix} \cdot \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \right|^2$$

$$= N^2 (1) = (0.6018)^2 \approx 0.36.$$

36% will come out on the top channel.

Q3.

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(a) Eq. (5) suggests \hat{H} has eigenstates $|1\rangle, |2\rangle$ and $|3\rangle$ with eigenvalues $\hbar\omega, 2\hbar\omega, 3\hbar\omega$ resp.

$$\text{So } |\psi(t)\rangle = e^{-i \frac{\hat{H} t}{\hbar}} \left(\frac{1}{\sqrt{2}} |1\rangle - \frac{i}{\sqrt{3}} |3\rangle \right)$$

$$= \frac{1}{\sqrt{2}} e^{-i\omega t} |1\rangle - \frac{i}{\sqrt{3}} e^{-3i\omega t} |3\rangle$$

(b) $\langle \hat{B} \rangle = \langle \psi(t) | \hat{B} | \psi(t) \rangle$

$$= \left(\frac{1}{\sqrt{2}} e^{+i\omega t} \langle 1| + \frac{i}{\sqrt{3}} e^{+3i\omega t} \right)$$

$$\times \left(2 |1\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2| \right)$$

$$\left(\frac{1}{\sqrt{2}} e^{-i\omega t} |1\rangle - \frac{i}{\sqrt{3}} e^{-3i\omega t} |3\rangle \right)$$

$$= b \left(\frac{1}{\sqrt{2}} e^{+i\omega t} |1\rangle + \frac{i}{\sqrt{3}} e^{+3i\omega t} \right)$$

$$\left(\frac{2}{\sqrt{2}} e^{-i\omega t} |1\rangle - \frac{i}{\sqrt{3}} e^{-3i\omega t} |2\rangle \right)$$

$$= b \left(\frac{2}{2} + \frac{1}{3} \right) = b \left(\frac{8}{6} \right) = \frac{4b}{3}$$

(c) At $t=0$, $|\psi(t=0)\rangle = |1\rangle$, an eigenstate of \hat{H} .

$$\hat{A} = a (|1\rangle\langle 2| + |2\rangle\langle 1| + 2|3\rangle\langle 3|)$$

In the eigenbasis of \hat{H} ,

$$\hat{A} = \begin{pmatrix} 0 & a & 0 \\ a & 0 & 0 \\ 0 & 0 & 2a \end{pmatrix}$$

Eigenstates and eigenvalues of \hat{A} are :

Eigenstates	Eigenvalues (measurement outcomes)
$\frac{1}{\sqrt{2}} (1\rangle + 2\rangle)$	a
$\frac{1}{\sqrt{2}} (1\rangle - 2\rangle)$	a
$ 3\rangle$	$2a$

Since outcome is $2a$, new state created is $|3\rangle$.

At time, t , the state will evolve into :

$$|\psi(t)\rangle = e^{-i\frac{\hat{A}t}{\hbar}} |3\rangle$$

$$= e^{-3i\omega t} |3\rangle.$$

Then \hat{B} is measured. We need to find outcomes and corresponding eigenstates.

$$\hat{B} = b (2 |1\rangle\langle 1| + |2\rangle\langle 3| + |3\rangle\langle 2|)$$

$$= \begin{pmatrix} 2b & 0 & 0 \\ 0 & 0 & b \\ 0 & b & 0 \end{pmatrix}$$

Eigenstates

Eigenvalues

$|1\rangle$

$2b$

$\frac{|2\rangle + |3\rangle}{\sqrt{2}}$

b

$\frac{|2\rangle - |3\rangle}{\sqrt{2}}$

b

Hence possible outcomes are ~~$|2\rangle$~~ , ~~$|3\rangle$~~

~~$|2\rangle$~~ ~~$|3\rangle$~~

$2b, b, b$. Let's find their

probabilities given the state
 $e^{-3i\omega t} |3\rangle = |3\rangle$ (ignore global phase).

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$$\text{Prob}(\text{output } 2b) = |\langle 1 | 3 \rangle|^2 = 0$$

$$\begin{aligned} \text{Prob}(\text{output } b) &= \left| \left(\frac{\langle 2 | + \langle 3 |}{\sqrt{2}} \right) | 3 \rangle \right|^2 \\ &+ \left| \left(\frac{\langle 2 | - \langle 3 |}{\sqrt{2}} \right) | 3 \rangle \right|^2 \\ &= \frac{1}{2} + \frac{1}{2} = 1 \end{aligned}$$