

Question 1

The time-dependent Schrödinger equation can be used to determine the evolution of a quantum state after some time t . It is written as:

$$i\hbar \frac{d|\psi(t)\rangle}{dt} = \hat{H} |\psi(0)\rangle,$$

where \hat{H} is the Hamiltonian of the system. In this question, we define \hat{H} to be:

$$\hat{H} = b\hbar\hat{X}.$$

- (a) Show that $\hat{X}^2 = \mathbf{1}$.
- (b) A quantum state $|\psi(0)\rangle$ at time $t = 0$ is given by

$$\frac{1}{\sqrt{2}} \left(|0\rangle + i|1\rangle \right).$$

What is $|\psi(t)\rangle$ if $bt = \frac{\pi}{2}$?

- (c) Identify $|\psi(0)\rangle$ and $|\psi(t)\rangle$ on the Bloch sphere.
- (d) If $bt = \frac{\pi}{2}$, what rotation will take $|\psi(0)\rangle$ to $|\psi(t)\rangle$? Specify the amount and axis of rotation.
- (e) Now suppose that the initial state is

$$\frac{1}{\sqrt{2}} \left(|0\rangle + |1\rangle \right).$$

How does this state evolve with time?

- (f) For part (f), is there any measurement that can differentiate between $|\psi(0)\rangle$ and the state at some later time?

The three Pauli operators are defined below

$$\begin{aligned} \hat{X} &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ \hat{Y} &= \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \\ \hat{Z} &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}. \end{aligned}$$

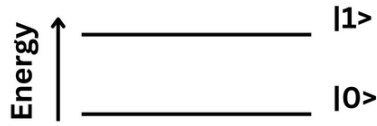


Figure 1: Energy Level Diagram for Question 2

Question 2

If an electron is in the lower energy state, it is labeled as $|0\rangle$, and if it is in the higher (excited) state, it is labeled as $|1\rangle$. A pulse is applied that has the Hamiltonian

$$\hbar V_0 \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

Our electron starts in the lower energy state. The pulse switches on at time $t = 0$ and ends at $\frac{T}{V_0}$.

What is the probability that the electron is found in the excited state at the end of the pulse?

Question 3

Show that a π rotation about the z -axis followed by a $\frac{\pi}{2}$ rotation about the y -axis generates a Hadamard gate.

Quiz 3 Group 1

Q1.

$$a) \hat{X} \hat{z} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \Rightarrow \hat{X}^2 \hat{z} = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$= \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \hat{z} = \hat{z}$$

since acting $\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$ on any $\begin{bmatrix} a \\ b \end{bmatrix}$ gives back $\begin{bmatrix} a \\ b \end{bmatrix}$.

$$b) |\psi(0)\rangle = \frac{1}{\sqrt{2}} (|0\rangle + i|1\rangle)$$

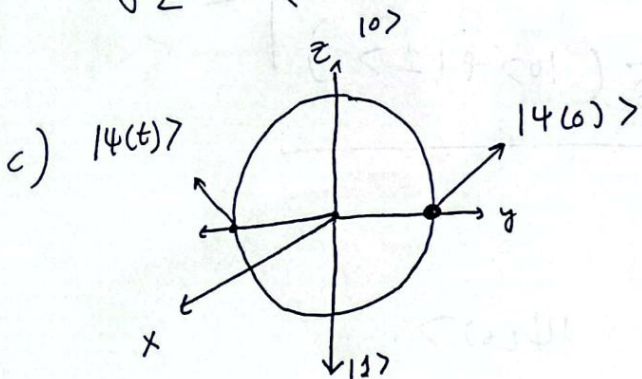
$$|\psi(t)\rangle = e^{-\frac{i}{\hbar} \hat{H} t} |\psi(0)\rangle = \frac{e^{-ibt\hat{X}}}{\sqrt{2}} (|0\rangle + i|1\rangle)$$

Now $e^{-ibt\hat{X}} = \cos(bt) \hat{1} - i \sin(bt) \hat{X}$

with $bt = \frac{\pi}{2}$, we are left with

$-i\hat{X}$ only. Acting this on $|\psi(0)\rangle$ yields

$$\frac{1}{\sqrt{2}} (-i|1\rangle + |0\rangle) = \boxed{\frac{1}{\sqrt{2}} (|0\rangle - i|1\rangle)}$$



$$\frac{|\psi(0)\rangle}{\theta = \frac{\pi}{2}, \phi = \frac{\pi}{2}}$$

$$\frac{|\psi(t)\rangle}{\theta = \frac{\pi}{2}, \phi = \frac{3\pi}{2}}$$

d) (π) . Angle $\rightarrow \pi$
 Axis $\rightarrow x$. (anticlockwise).

e) $|\psi(t)\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$.

$$|\psi(t)\rangle = (\cos(bt) \mathbb{1} - i \sin(bt) \hat{X}) \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$= \frac{1}{\sqrt{2}} [(\cos(bt) - i \sin(bt)) |0\rangle + (\cos(bt) - i \sin(bt)) |1\rangle]$$

$$= \frac{1}{\sqrt{2}} e^{-ibt} (|0\rangle + |1\rangle) = \frac{e^{-ibt}}{\sqrt{2}} (|0\rangle + |1\rangle) = |\psi(t)\rangle$$

f) No. Suppose we measure the probability of getting some state as the outcome, $|n\rangle$.

for $|\psi(t)\rangle =$

$$\left| \langle n | \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \right|^2$$

for $|\psi(t)\rangle =$

$$\left| \langle n | \frac{e^{-ibt}}{\sqrt{2}} (|0\rangle + |1\rangle) \right|^2$$

$$= |e^{-ibt}|^2 \left| \langle n | \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \right|^2$$

$$= 1$$

\Rightarrow same as for $|\psi(t)\rangle$.

Q2.

$$\hat{H} = \hbar V_0 \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \hbar V_0 \hat{X}$$

$$|4(0)\rangle = |0\rangle$$

$$|4(t)\rangle = e^{-\frac{i}{\hbar} (\hbar V_0 \hat{X}) t} |0\rangle$$

$$= e^{-i V_0 t \hat{X}} |0\rangle$$

We are interested at $t = \frac{T}{V_0}$.

$$|4\left(\frac{T}{V_0}\right)\rangle = e^{-iT \hat{X}} |0\rangle$$

$$= \left[\cos(T) \mathbb{1} - i \sin(T) \hat{X} \right] |0\rangle$$

$$= \cos(T) |0\rangle - i \sin(T) |1\rangle$$

$$\text{Prob}(|1\rangle) =$$

$$|\langle 1 | \psi(t) \rangle|^2 = |-i \sin(T)|^2$$

$$= (-i \sin(T)) (i \sin(T))$$

$$= \sin^2(T)$$

Q3.

$$|0\rangle \xrightarrow{\left(\frac{\pi}{2}\right)_z} e^{-\frac{i\pi}{2} \hat{Z}} |0\rangle = \left[\cos\left(\frac{\pi}{2}\right) \mathbb{1} - i \sin\left(\frac{\pi}{2}\right) \hat{Z} \right] |0\rangle$$

$$= -i |0\rangle$$

$$-i |0\rangle \xrightarrow{\left(\frac{\pi}{2}\right)_y} e^{-\frac{i\pi}{4} \hat{Y}} (-i |0\rangle) = \frac{-i}{\sqrt{2}} (\mathbb{1} - i \hat{Y}) |0\rangle$$

$$= \frac{-i}{\sqrt{2}} [|0\rangle + |1\rangle]$$

$$|1\rangle \xrightarrow{(\pi)_z} i|1\rangle$$

Same steps as before

but note that $\hat{z}|0\rangle = |0\rangle$ but

$$\hat{z}|1\rangle = -|1\rangle$$

$$i|1\rangle \xrightarrow{(\frac{\pi}{2})_y} e^{-\frac{i\pi}{4}} i|1\rangle$$

$$= \frac{1}{\sqrt{2}} [10\rangle - i(11\rangle)]$$

$$= \frac{1}{\sqrt{2}} (11\rangle - 10\rangle)$$

$$= \frac{i}{\sqrt{2}} (11\rangle - 10\rangle)$$

$$= \frac{-i}{\sqrt{2}} (10\rangle - 11\rangle)$$

we could not have said what we say here if the phases for $|0\rangle, |1\rangle$ were different!

Since we pick up the same global phase with $|0\rangle \hat{z} |1\rangle$, so our rotations will perform the Hadamard gate on any arbitrary state. It will give the correct state with a global phase multiplied. But this can be dropped.

Note that $-i = e^{i\frac{3\pi}{2}} = e^{-i\frac{\pi}{2}}$.

Also any arbitrary state is $a|0\rangle + b|1\rangle$ and the Hadamard gate acts on each term by turn.

